Transverse foliations in the Euler problem of two fixed centers

Geometry, Dynamics and Mechanics Seminar

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The Euler problem of two centers in the plane



In canonical coordinates:

$$H_{\mu} = \frac{p_x^2 + p_y^2}{2} - \frac{\mu}{\sqrt{(x+1)^2 + y^2}} - \frac{1-\mu}{\sqrt{(x-1)^2 + y^2}}$$

• mass ratio: $0 < \mu < 1$.

Euler (1760): integral of motion

$$G = \frac{1}{2}(xp_y - yp_x)^2 + \frac{p_x^2}{2} + x\frac{\mu}{\sqrt{(x+1)^2 + y^2}} - x\frac{1-\mu}{\sqrt{(x-1)^2 + y^2}}$$

Hill's regions



$$egin{aligned} & H_\mu = c \Rightarrow V(x,y) \leq c \ & c_{ ext{crit}} = -rac{1}{2} - \sqrt{\mu - \mu^2}. \end{aligned}$$

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- $c < c_{crit}$: trajectories stay close to a center.
- $c_{\text{crit}} < c < 0$: transit trajectories.
- c > 0: the typical trajectory is unbounded.

Some related models in Celestial Mechanics

circular planar restricted three-body problem

$$H_{\mu} = \frac{p_x^2 + p_y^2}{2} + (x - \mu)p_x - yp_x - \frac{\mu}{\sqrt{(x - 1)^2 + y^2}} - \frac{1 - \mu}{\sqrt{x^2 + y^2}}$$

Hill's lunar problem

$$H = \frac{p_x^2 + p_y^2}{2} + yp_x - xp_y - \frac{1}{\sqrt{x^2 + y^2}} - x^2 + \frac{y^2}{2}.$$

Rotating Kepler problem

$$H = \frac{p_x^2 + p_y^2}{2} + yp_x - xp_y - \frac{1}{\sqrt{x^2 + y^2}}.$$

Harmonic oscillator

$$H = \frac{p_x^2 + p_y^2}{2} + \frac{x^2 + y^2}{2}.$$

Hénon-Heiles system

$$H = \frac{p_x^2 + p_y^2}{2} + \frac{x^2 + y^2}{2} + x^2 y - \frac{y^3}{3}.$$

Reeb flows

- ► *M* compact, connected 3-manifold.
- \blacktriangleright λ contact form on M:

 $\lambda \wedge d\lambda \neq 0.$

- $\xi = \ker \lambda \subset TM$ contact structure.
- \blacktriangleright Reeb vector field R

$$d\lambda(R,\cdot)=0$$
 and $\lambda(R)=1.$

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The Reeb flow parametrizes Hamiltonian flows on contact-type energy levels.

Transverse foliations for Reeb flows

A transverse foliations consists of a singular foliation of ${\cal M}$ adapted to the Reeb flow:

- A link $L \subset M$ formed by finitely many closed orbits.
- ► a smooth foliation of M \ L by properly embedded surfaces transverse to the flow. The closure of each leaf is a compact surface with boundary contained in L.



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Why transverse foliations?

- some order inside chaos.
- two-dimensional reduction.
- existence and multiplicity of periodic orbits.

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- symbolic dynamics.
- the dynamics of transit trajectories.
- homoclinics to the Lyapunoff orbit.

Pseudo-holomorphic curves

- Let $J: \xi \to \xi$ be a $d\lambda$ -compatible complex structure.
- Consider the ℝ-invariant almost complex structure J̃ on the symplectization ℝ × M so that:

$$\widetilde{J}|_{\xi} = J$$
 and $\widetilde{J} \cdot \partial_a = R$

where ∂_a is the unit vector in the \mathbb{R} -direction. $\tilde{u} = (a, u) : (\dot{\Sigma}, i) \to (\mathbb{R} \times M, \widetilde{J})$ is pseudo-holomorphic:

$$d\tilde{u}\circ i=\widetilde{J}(u)\circ d\tilde{u}.$$

Finite Hofer's energy

$$0 < E(\tilde{u}) := \sup_{\varphi} \int_{\dot{\Sigma}} \tilde{u}^* d(\varphi(a)\lambda) < +\infty,$$

where $\varphi : \mathbb{R} \to [0,1]$ satisfies $\varphi' \ge 0$.

Periodic orbits at the ends of a finite energy curve



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Why pseudo-holomorphic curves?

- pseudo-holomorphic curves can be found almost explicitly for many integrable Reeb flows.
- one can deform the pseudo-holomorphic curves to find periodic orbits, global surfaces of section and transverse foliations for Reeb flows far from integrable ones.

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Symplectic cobordism



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Hofer-Wysocki-Zehnder (Ann. of Math. 1998)

Back to the Euler problem: elliptic coordinates



New canonical coordinates: (u, v, p_u, p_v) :

 $x = \cosh u \cos v$ $y = \sinh u \sin v$ $p_x dx + p_y dy = p_u du + p_v dv$

Regularized Hamiltonian:

$$K_{\mu,c} = (\cosh^2 u - \cos^2 v)(H_{\mu} - c)$$

= $\frac{p_v^2}{2} + (2\mu - 1)\cos v + c\cos^2 v + \frac{p_u^2}{2} - \cosh u - c\cosh^2 u$

Energy x mass ratio



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▶ II, III and IV - $c_{crit} < c < 0$ (high energy). ▶ I - $c < c_{crit} = -\frac{1}{2} - \sqrt{\mu - \mu^2}$ (low energy).

Main theorem

Theorem(de Paulo, Hryniewicz, Kim) For every $0 < \mu < 1$ and $c_{\rm crit} < c < 0$, the regularized energy surface $K_{\mu,c}(0)$ admits a transverse foliation, which is the projection of a finite energy foliation. The binding orbits are the Lyapunoff orbit and the exterior collision orbits.



Low energy *I*: $c < c_{crit}(\mu)$



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- ▶ 2 components diffeomorphic to S^3 .
- disk-like global surface of section

Low energy *I*: $c < c_{crit}(\mu)$



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Open book decomposition.

Each page is a disk-like global surface of section.

High energy $II : c_{crit} < c < -\frac{1}{2}$



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▶ 1 component diffeomorphic to $S^1 \times S^2$.

High energy $II : c_{crit} < c < -\frac{1}{2}$



Weakly convex foliation: 2 pairs of rigid cylinders, 2 pairs of rigid planes and 2 families of planes.

Pseudo-holomorphic

- some planes and cylinders can be made projections of pseudo-holomorphic curves.
- take action-angle coordinates in the (u, p_u) plane.
- choose a suitable decoupled contact form on $K_{\mu,c}^{-1}(0)$.
- use the complex structure $J: \xi \to \xi$ induced by the quaternions.
- consider solutions of the form

$$(u, p_u, v, p_v) = (r(s) \cos 2\pi t, r(s) \sin 2\pi t, 0, g(s)),$$

 $\forall (s,t) \in \mathbb{R} \times (\mathbb{R}/\mathbb{Z}).$

solve first order ODEs for r(s) and f(s) to obtain a pseudo-holomorphic plane.

Symbolic dynamics



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- ► Satellite: 1111111... or 2222222...
- Lemniscate: 34343434343434.....

High energy $III : -\frac{1}{2} < c < c_2(\mu)$



▶ 1 component diffeomorphic to $S^1 \times S^2$.

High energy $III : -\frac{1}{2} < c < c_2(\mu)$



Weakly convex foliation: 2 pairs of rigid cylinders, 2 pairs of rigid planes and 2 families of planes.

Symbolic dynamics



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- Planet: 343434343434....

High energy $IV : c_2(\mu) < c < 0$



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▶ 1 component diffeomorphic to $S^1 \times S^2$.

High energy IV : $c_2(\mu) < c < 0$



Weakly convex foliation: 1 pair of rigid cylinders, 1 pair of rigid planes and 1 family of planes.

Symbolic dynamics



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Hénon-Heiles system



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