## Analysis on manifolds and applications to layer potentials

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## Dedicated to the memory of Gabriela Kohr

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## Introduction

These are the lecture notes that I will use for my five-hour mini-course at the conference. I plan to do "chalk and board" talks in the beginning, so I will not use these slides for all the lectures.

These notes contain more material then what I intend to present in my mini-course. In particular, the last section (the second part of the lecture notes) was not covered in the mini-course. (This second part contains an attempt to generalize the results of first part.)

At this time, these notes have not been revised carefully. In particular, not enough credit is given to other researchers, and I sincerely appologize for this omission. However, I include a very short (and extremely incomplete) reference section at the end of each of the two parts of the lecture notes.

I have two main goals for my mini-course:

- First, to give a very quick introduction to some of the main tools and concepts used in analysis on manifolds.
- Second, to use some of these tools to investigate the method of layer potentials on manifolds with cylindrical ends.
Even if one is interested in domains in $R^{n}$, the method of layer potentials require us to work on a non-trivial manifold, namely the boundary of our domain.
- If time permits, I will state the well-posedness of the Poisson problem on manifolds with finite width.

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## Motivation

- Even if one is interested in analysis on Euclidean spaces, one is lead to consider (smooth) manifolds, because they are the boundaries of smooth domains.
- Because certain domains are not smooth (i.e. do not have a smooth boundary), we are also lead to consider singular spaces. Example: domains with conical points (including polygons).
- A related example: domains with conical points. (Our method to study domains with conical points is to use domains with cylindrical ends.)
- More complicated domains, such as polyhedral domains, require the use of Lie manifolds, discussed in the last section.
- Very many researchers: Stein, Strichartz, Browder, ... (early works on analysis on non-compact manifolds), Dauge, Kondratiey (eatis Melrose, Schulze, ...


## Smooth manifolds

## Definition

A smooth manifold $M$ is a locally Euclidean, second countable Hausdorff space endowed with a smooth structure.

Thus $M=\cup_{\alpha \in I} U_{\alpha}$ is Hausdorff, second countable, with each $U_{\alpha}$ an open subset of $M$ and there exist homeomorphisms $\phi_{\alpha}: U_{\alpha} \rightarrow W_{\alpha} \subset \mathbb{R}^{n}$, such that, for all $\alpha, \beta \in I$, the induced map

$$
\phi_{\alpha \beta}:=\phi_{\alpha} \circ \phi_{\beta}^{-1}: \phi_{\beta}\left(U_{\alpha} \cap U_{\beta}\right) \rightarrow \phi_{\beta}\left(U_{\alpha} \cap U_{\beta}\right)
$$

## is a diffeomorphism.

A pair $\left(U_{\alpha}, \phi_{\alpha}\right)$ is called a coordinate chart and the set $\left\{U_{\alpha} \mid \alpha \in I\right\}$ is called an atlas on $M$.
$f: M \rightarrow \mathbb{R}$ is smooth (i.e. in $\mathcal{C}^{\infty}(M)$ ) if all $f \circ \phi_{\alpha}^{-1}: W_{\alpha} \rightarrow \mathbb{R}$ are smooth. $M$ will always denote a smooth manifold.

## Remark on the definition of smooth manifolds

Some authors replace second countability by the weaker requirement that the manifold be a paracompact topological space. The second choice implies that every connected component is second countable.

Thus, a manifold in the second sense is a manifold in the first sense if, and only if, the set of connected components is countable. For instance, an uncountable set $D$ with the discrete topology is not a manifold in the first sense, but it is a 0 -dimensional manifold in the second sense.

For our results (and for most statements in differential geometry), it does not matter which definition we choose (second countable or paracompact).

## Vector bundles

Let $E, X$ be topological vector space together with a map $\pi: E \rightarrow X$. For every $A \subset X$, we shall write $\left.E\right|_{A}:=\pi^{-1}(A) \subset E$.

## Definition

$(E, \pi)$ is a (real, topological) vector bundle over $X$ if

- For every $p \in X$, the set $E_{p}:=\pi^{-1}(p)$ is a real vector space;
- there exists an open covering $M=\cup_{\alpha \in I} \cup_{\alpha}$ and a fiber-preserving* homeomorphisms

$$
\psi_{\alpha}:\left.E\right|_{U_{\alpha}} \rightarrow U_{\alpha} \times \mathbb{R}^{n_{\alpha}}
$$

that is linear in each fiber.

* (i.e. $\psi_{\alpha}\left(E_{q}\right)=\{q\} \times \mathbb{R}^{n_{\alpha}}$, equivalently, $\pi_{1}\left(\psi_{\alpha}(x)\right)=\pi(x)$.)


## Consequences

It follows that, for any two $\alpha, \beta \in X$, the resulting map

$$
\psi_{\alpha \beta}:=\psi_{\alpha} \circ \psi_{\beta}^{-1}:\left(U_{\alpha} \cap U_{\beta}\right) \times \mathbb{R}^{n_{\beta}} \rightarrow\left(U_{\alpha} \cap U_{\beta}\right) \times \mathbb{R}^{n_{\alpha}}
$$

is a homeomorphism linear in each fiber, thus of the form

$$
\psi_{\alpha \beta}(x, v)=\left(x, A_{\alpha \beta}(x) v\right)
$$

where

$$
A_{\alpha \beta}: U_{\alpha} \cap U_{\beta} \rightarrow G L\left(\mathbb{R}^{n_{\beta}}\right):=\operatorname{End}\left(\mathbb{R}^{n_{\beta}}\right)^{-1}
$$

is continuous (the famous "transition functions").
Alternatively, the continuous vector bundle can be defined starting from the transition functions, as long as they satisfy the cocycle condition:

$$
A_{\alpha \beta}(x) A_{\beta \gamma}(x)=A_{\alpha \gamma}(x), \quad \text { for } x \in U_{\alpha} \cap U_{\beta} \cap U_{\gamma}
$$

## Smooth vector bundles

Using the cocycle condition, we can identify

$$
E \simeq \bigsqcup\left(U_{\alpha} \times \mathbb{R}^{n_{\alpha}}\right) / \sim
$$

(disjoint union), where $(x, v) \sim\left(x, A_{\alpha \beta}(x) v\right)$ (with the first $x$ in $U_{\beta}$ and the second one in $U_{\alpha}$ and $v \in \mathbb{R}^{n_{\beta}}$ ).

Our vector bundles are thus locally trivial.
The simplest example of a vector bundle is the trivial vector bundle

$$
E:=X \times \mathbb{R}^{n} .
$$

Then all $U_{\alpha}=X$ and the transition functions are the identity matrices.
Complex vector bundles are treated similarly.

## Continuous sections of a vector bundle

## Definition

Let $(E, \pi)$ be a topological vector bundle over $X$. A (continuous) section of $E$ is a continuous map $s: X \rightarrow E$ such that

$$
\pi \circ s=i d
$$

The last condition means that $s(x) \in E_{x}:=\pi^{-1}(x)$, for all $x \in X$.
If $E=X \times \mathbb{R}^{N}$ (trivial of rank $N$ ), then a section $s: X \rightarrow E$ is simply an $n$-uple of continuous functions:

$$
s(x)=\left(x, s_{1}(x), s_{2}(x), \ldots, s_{n}(x)\right)
$$

$s_{j}: X \rightarrow \mathbb{R}$ continuous.

## Smooth vector bundles

## Definition

Let $(E, \pi)$ be a vector bundle over $X$. Let us assume that $X$ is a smooth manifold and that $\pi$ and the transition functions $A_{\alpha \beta}$ are smooth. Then $E$ is a smooth vector bundle (and a smooth manifold).

If $U \subset \mathbb{R}^{n}$ is an open subset, then it is a smooth manifold and

$$
T U \simeq U \times \mathbb{R}^{n}
$$

In general, if $M$ is a smooth manifold with coordinate charts $\left(U_{\alpha}, \phi_{\alpha}\right)$, then the differentials

$$
D \phi_{\alpha \beta}(x) \in M_{n}(\mathbb{R})^{-1}
$$

satisfy the cocycle condition and can be used to define a smooth (real) vector bundle $\pi: T M \rightarrow M$, called the tangent bundle of $M$.

## The tangent bundle

Let $M$ be a smooth manifold. Recall that a function $f: M \rightarrow \mathbb{R}$ is in $\mathcal{C}^{\infty}(M)$, iff $f \circ \phi_{\alpha}^{-1}$ is smooth on $W_{\alpha}:=\phi_{\alpha}\left(U_{\alpha}\right)$.

An element $v \in T_{p} M=(T M)_{p}=\pi^{-1}(p)$ can be defined intrinsically as a derivation of $\mathcal{C}^{\infty}(M)$ at $p$, that is, a linear map $v: \mathcal{C}^{\infty}(M) \rightarrow \mathbb{R}$ s.t.

$$
v(f g)=v(f) g(p)+f(p) v(g)
$$

Similarly, a smooth section $X$ of $T M$ can be identified with a derivation $X: \mathcal{C}^{\infty}(M) \rightarrow \mathcal{C}^{\infty}(M)$. For instance, if $U \subset \mathbb{R}^{n}$ is an open subset, a section $X$ of $T U=U \times \mathbb{R}^{n}$ is of the form $X=\sum_{k=1}^{n} X_{k} \frac{\partial}{\partial x_{k}}$, explicitly

$$
X u(x)=\sum_{k=1}^{n} X_{k}(x) \frac{\partial u}{\partial x_{k}}(x)
$$

Moreover, $[X, Y] u:=X(Y u)-Y(X u)$ Lie bracket is also a derivation (i.e. a vector field.)

## Covariant derivatives

For reasons that will become apparent in applications, we want to work not just with smooth functions on $M$, but also with smooth sections $\mathcal{C}^{\infty}(E)=\mathcal{C}^{\infty}(M ; E)$ of various smooth vector bundles $E \rightarrow M$.
Note that $\mathcal{C}^{\infty}(E)$ is a $\mathcal{C}^{\infty}(M)$-module.
The big question that we face then is how to differentiate the smooth sections of such a bundle (not a unique solution).

## Definition

A covariant derivative $\nabla=\nabla^{E}$ on $E \rightarrow M$ is a bilinear map $\nabla=\nabla^{E}: \mathcal{C}^{\infty}(T M) \times \mathcal{C}^{\infty}(E) \rightarrow \mathcal{C}^{\infty}(E)$ s.t.
(1) $\nabla_{f x}(\xi)=f \nabla_{x}(\xi)$
(2) $\nabla_{X}(f \xi)=X(f) \xi+f \nabla_{X} \xi$ (Leibnitz's rule),
$\forall \xi \in \mathcal{C}^{\infty}(E), X \in \mathcal{C}^{\infty}(T M), f \in \mathcal{C}^{\infty}(M)$.

## The local form of covariant derivatives I

Let us assume that $U \subset \mathbb{R}^{n}$ is an open subset and that $E=U \times \mathbb{R}^{N}$ is a trivial vector bundle.

A possible choice of a connection $\nabla: \mathcal{C}^{\infty}(T U) \times \mathcal{C}^{\infty}(E) \rightarrow \mathcal{C}^{\infty}(E)$ is

$$
\nabla_{x} u=X u:=\sum_{k=1}^{n} x_{k} \frac{\partial}{\partial x_{k}} u,
$$

where $X=\sum_{k=1}^{n} x_{k} \frac{\partial}{\partial x_{k}}$ and the derivation is component-wise. That is

$$
\nabla_{\frac{\partial}{\partial x_{k}}} u=\left(\frac{\partial}{\partial x_{k}} u_{1}, \frac{\partial}{\partial x_{k}} u_{2}, \ldots, \frac{\partial}{\partial x_{k}} u_{n}\right),
$$

where $u=\left(u_{1}, u_{2}, \ldots, u_{n}\right)$. This is the trivial connection on the trivial vector bundle $E=U \times \mathbb{R}^{N}$.

## The local form of covariant derivatives II

The choice

$$
\nabla_{X} u=X u:=\sum_{k=1}^{n} X_{k} \frac{\partial}{\partial x_{k}} u
$$

for the connection $\nabla: \mathcal{C}^{\infty}(T U) \times \mathcal{C}^{\infty}(E) \rightarrow \mathcal{C}^{\infty}(E)$ on the trivial bundle $E=U \times \mathbb{R}^{N}$ is, however, not unique, since, given any choice of smooth matrix functions $A_{k} \in \mathcal{C}^{\infty}\left(M ; M_{N}\right)$, the formula

$$
\nabla_{X}^{\prime} u=X u+\left(\sum_{k=1}^{n} x_{k} A_{k}\right) u:=\sum_{k=1}^{n} x_{k}\left(\frac{\partial}{\partial x_{k}} u+A_{k} u\right)
$$

also defines a connection on $E \rightarrow U$.
These are all connections on $E$.

## The local form of covariant derivatives III

We thus have

$$
\nabla_{X}^{\prime} u-\nabla_{X} u=\left(\sum_{k=1}^{n} X_{k} A_{k}\right) u=: \mathcal{A}(X) u
$$

In general, i.e., given two covariant derivatives $\nabla$ and $\nabla^{\prime}$ on any $E \rightarrow M$

$$
\mathcal{A}(X):=\nabla_{X}^{\prime}-\nabla_{X}=\mathcal{C}^{\infty}(M ; \operatorname{End}(E)),
$$

We can regard then $\mathcal{A} \in \Omega^{1}(M ; \operatorname{End}(E)):=\mathcal{C}^{\infty}\left(M ; T^{*} M \otimes \operatorname{End}(E)\right)$ and $\nabla=\nabla^{E}$ as a differential operator

$$
\nabla: \mathcal{C}^{\infty}(E) \rightarrow \mathcal{C}^{\infty}\left(T^{*} M \otimes E\right)
$$

satisfying $\nabla_{X} u=\langle X, \nabla u\rangle$ and $\nabla(f \xi)=d f \otimes \xi+f \nabla \xi$.

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## Operations with vector bundles

If $E, F \rightarrow M$ are two (smooth) vector bundles, we can define the (smooth) vector bundles

$$
E^{*}, \operatorname{End}(E), \operatorname{Hom}(E ; F), E \otimes F, \Lambda^{\kappa} E \rightarrow M,
$$

by defining their fibers
$\left(E^{*}\right)_{x}:=\left(E_{x}\right)^{*}, \operatorname{End}(E)_{x}:=\operatorname{End}\left(E_{x}\right), \operatorname{Hom}(E ; F)_{x}:=\operatorname{Hom}\left(E_{x} ; F_{x}\right), \ldots$
for any $x \in M$. Their topology (or smooth structure) is obtained from the cocycle construction.

Connections on $E$ and $F$ will define natural connections on these new (derived) bundles.
A metric $g$ on $E$ is a smooth section of $E^{*} \otimes E^{*} \simeq(E \otimes E)^{*}$ that is a metric on each fiber.

## Connections on the new bundles

Recall: connection $\nabla=\nabla^{E}$ on $E \rightarrow M$ is a differential operator

$$
\nabla: \mathcal{C}^{\infty}(E) \rightarrow \mathcal{C}^{\infty}\left(T^{*} M \otimes E\right) \text { s.t. } \nabla_{X} u=\langle X, \nabla u\rangle .
$$

- Connection on $E^{*}$ : For $\xi \in \mathcal{C}^{\infty}\left(E^{*}\right)$ and $\eta \in \mathcal{C}^{\infty}(E)$, we let

$$
\left\langle\nabla_{x} \xi, \eta\right\rangle:=X\langle\xi, \eta\rangle-\left\langle\xi, \nabla_{x} \eta\right\rangle .
$$

- Connection on $E \otimes F$ : we consider the connection

$$
\nabla^{E \otimes F}:=\nabla^{E} \otimes 1+1 \otimes \nabla^{F},
$$

that is

$$
\nabla_{X}^{E \otimes F}(\xi \otimes \eta):=\nabla_{X}^{E} \xi \otimes \eta+\xi \otimes \nabla_{X}^{F} \eta .
$$

## Higher order covariant derivatives

We endow TM with the Levi-Civita connection

$$
\nabla^{M}: \mathcal{C}^{\infty}(M ; T M) \rightarrow \mathcal{C}^{\infty}\left(M ; T^{*} M \otimes T M\right)
$$

This allows us to iterate:

$$
\begin{aligned}
\nabla^{k}: \mathcal{C}^{\infty}(E) \xrightarrow{\nabla} \mathcal{C}^{\infty}\left(T^{*} M \otimes E\right) \xrightarrow{\nabla} \mathcal{C}^{\infty}\left(T^{*} M \otimes\right. & \left.T^{*} M \otimes E\right) \\
& \xrightarrow{\nabla} \mathcal{C}^{\infty}\left(T^{* \otimes k} M \otimes E\right),
\end{aligned}
$$

where endow each tensor bundle $T^{* \otimes k} M \otimes E$ with the induced (Levi-Civita) tensor product connection:

$$
\begin{aligned}
& \nabla_{X}\left(\xi_{1} \otimes \ldots \otimes \xi_{k} \otimes \eta\right):=\nabla_{X}^{M}\left(\xi_{1}\right) \otimes \ldots \otimes \eta+\xi_{1} \otimes \nabla_{X}^{M}\left(\xi_{2}\right) \otimes \ldots \otimes \eta \\
&+\ldots+\xi_{1} \otimes \xi_{2} \otimes \ldots \otimes \xi_{k} \otimes \nabla_{X}^{E}(\eta)
\end{aligned}
$$

## Metric preserving connections

Recall that a metric $g$ on $E$ is a smooth section of $E^{*} \otimes E^{*} \simeq(E \otimes E)^{*}$ that is a metric on each fiber. That is, $g$ is a smoothly varying family of metrics $g_{x}$ on $E_{x}:=\pi^{-1}(x)$.

The connection $\nabla$ is called metric preserving if, $\nabla g=0$, explicitly

$$
g\left(\nabla_{x} \xi, \eta\right)+g\left(\xi, \nabla_{X} \eta\right)=X g(\xi, \eta)
$$

A metric $g$ on $M$ is a metric on $T M$, in which case $M$ is called Riemannian.
$(M, g)$ will be a Riemannian manifold from now on.UNIVERSITEE
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## The Levi-Civita connection

The Levi-Civita connection $\nabla^{M}$ may be defined using Koszul's formula: If $X, Y, Z$ are vector fields (i.e. in $\mathcal{C}^{\infty}(M ; T M)$ ), then

$$
\begin{aligned}
2 g\left(\nabla_{X}^{M} Y, Z\right)=X & (g(Y, Z))+Y(g(X, Z))-Z(g(X, Y)) \\
& -g([Y, X], Z)-g([X, Z], Y)-g([Y, Z], X)
\end{aligned}
$$

It is the unique metric preserving connection on TM that is also torsion-free, in the sense that

$$
\nabla_{X}^{M} Y-\nabla_{Y}^{M} X=[X, Y]
$$

for all (smooth) vector fields $X, Y \in \mathcal{C}^{\infty}(M ; T M)$.)

## Sobolev spaces

Let $E \rightarrow M$ be a finite dimensional vector bundle with metric and metric preserving connection $\nabla=\nabla^{E}$. We can define then $L^{p}(M ; E)$ with norm $\left(\int_{M}\|s\| d x\right)^{1 / p}$ (ess-sup if $p=\infty$ ).

## Definition

Let $k \in \mathbb{Z}_{+}$and $p \in[1, \infty]$. Then

$$
W^{k, p}(M ; E):=\left\{u \in L^{p}(M ; E) \mid \nabla^{j} u \in L^{p}\left(M ; T^{* \otimes j} M \otimes E\right), 1 \leq j \leq k\right\} / \equiv
$$ is the order $k, L^{p}$-type covariant Sobolev space of sections of $E$.

## Newtonian potential

The goal is to solve the Dirichlet boundary value problem

$$
\left\{\begin{array}{cl}
\Delta u=0 & \text { in } \Omega \\
u=f & \text { on } \partial \Omega
\end{array}\right.
$$

To explain the layer potentials method, recall that the solution of

$$
\Delta v=h \in \mathcal{C}_{C}\left(\mathbb{R}^{n}\right)
$$

is given by the Newtonian potential (fundamental solution of $\Delta$ ):

$$
v(x)=\Delta^{-1} h(x)=c_{n} \int_{\mathbb{R}^{n}}|x-y|^{2-n} h(y) d y
$$

## The method of layer potentials ( $\mathcal{S}$ )

The method of the single layer potential (for $\Delta$ ) is to try to represent $u$ in the form

$$
u(x):=\Delta^{-1}\left(g \otimes \delta_{\partial \Omega}\right)=\int_{\partial \Omega}|x-y|^{2-n} g(y) d \sigma
$$

More precisely, let:

- $d \sigma$ is the surface measure on $\partial \Omega$.
- $\left\langle h, g \otimes \delta_{\partial \Omega}\right\rangle:=\int_{\partial \Omega} h(x) g(x) d \sigma(x)$.

Then the single layer potential operator $\mathcal{S}$

$$
\mathcal{S} g(x):=\Delta^{-1}\left(g \otimes \delta_{\partial \Omega}\right)=\int_{\partial \Omega}|x-y|^{2-n} g(y) d \sigma
$$

## Harmonic function and boundary condition

The first formula:

$$
\mathcal{S} g(x):=\Delta^{-1}\left(g \otimes \delta_{\partial \Omega}\right)
$$

gives that $\mathcal{S g}$ is harmonic inside $\Omega$ :

$$
\Delta \mathcal{S} g(x)=\Delta\left(\Delta^{-1}\left(g \otimes \delta_{\partial \Omega}\right)\right)(x)=g \otimes \delta_{\partial \Omega}(x)=0
$$

## The second formula:

$$
\mathcal{S} g(x):=c_{n} \int_{\partial \Omega}|x-y|^{2-n} g(y) d \sigma(y), \quad x \in \Omega
$$

for $x \rightarrow z \in \partial \Omega$ gives directly

$$
\mathrm{S} g(z):=\left.\mathcal{S} g\right|_{\partial \Omega}(z):=c_{n} \int_{\partial \Omega}|z-y|^{2-n} g(y) d \sigma(y), \quad z \in \partial \Omega
$$

## Boundary conditions

Let (as on the previous slide)

$$
\mathrm{S} g(z):=\left.\mathcal{S} g\right|_{\partial \Omega}(z):=c_{n} \int_{\partial \Omega}|z-y|^{2-n} g(y) d \sigma(y), \quad z \in \partial \Omega
$$

(a "jump relation.")

Summary of the single layer potential method
Let $g=\mathbf{S}^{-1} f$. Then $u:=\mathcal{S} g$ solves the Dirichlet problem

$$
\begin{cases}\Delta u=0 & \text { in } \Omega \\ u=f & \text { on } \partial \Omega\end{cases}
$$

So we it is enough to invert $S$.

## The second lecture

## Double layer potential

We prefer the analogous double layer potential operator $\mathcal{D}$, namely:

$$
\mathcal{D} h(x):=c_{n} \int_{\partial \Omega} \partial_{\nu_{y}}|x-y|^{2-n} h(y) d \sigma(y)=\Delta^{-1}\left(f \otimes \delta_{\partial \Omega}^{\prime}\right)
$$

Difference: if we pass to the limit $x \rightarrow z \in \partial \Omega$, we obtain a "jump relation" that has some additional terms:

$$
\begin{aligned}
\lim _{x \rightarrow z} \mathcal{D} h(x)= & c_{n} \text { p.v. } \int_{\partial \Omega} \partial_{\nu_{y}}|z-y|^{2-n} h(y) d \sigma(y)+\frac{h(z)}{2} \\
& =\left(\mathbf{K}+\frac{1}{2}\right) h(z), \quad z \in \partial \Omega
\end{aligned}
$$

Summary of the double layer potential method
Our Dirichlet problem ( $\Delta u=0$ and $u=f$ on $\partial \Omega$ ) is reduced to

$$
\left(\frac{1}{2}+\mathbf{K}\right) h=f .
$$

## Pseudodifferential operators

Recall that (up to a factor) K is the integral operator with kernel

$$
\partial_{\nu y}|z-y|^{2-n}=(2-n) \frac{(z-y) \cdot \nu_{y}}{|z-y|^{n}} .
$$

Hence, if $\partial \Omega$ is smooth, then K is an order -1 pseudodifferential operator. (A similar fact holds for S , but will not be used.)

An order - 1 pseudodifferential operator on a compact manifold is compact. Thus

## Basic fact

If $\Omega$ is smooth and bounded, then K is compact.

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## Fredholm operators

## Definition

Let $T: X \rightarrow Y$ be a continuous, linear map between Banach spaces.
We say that $T$ is Fredholm if $T^{-1}(0)$ and $Y / T X$ are finite dimensional. Then its index is $\operatorname{ind}(T):=\operatorname{dim} \operatorname{ker}(T)-\operatorname{dim}(Y / T X)$.

## Easy observation

## An often used argument

If $T$ is Fredholm of index zero and injective, then it is an isomorphism (invertible).

We obtain that $\frac{1}{2}+\mathrm{K}$ is a Fredholm operator.
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## Well-posedness of the Dirichlet problem

We can now prove the Hadamard well posedness (solvability) of our Dirichlet problem ( $\Delta u=0$ and $\left.u\right|_{\partial \Omega}=f$ ) if $\Omega$ is smooth and bounded.

Recall that it is enough to invert $\frac{1}{2}+\mathrm{K}$, since our solution is

$$
u=\mathcal{D}\left(\left(\frac{1}{2}+K\right)^{-1} f\right)
$$

The proof of the invertibility of $\frac{1}{2}+\mathrm{K}$ is as follows:

## Invertibility of $\frac{1}{2}+K$ and solvability for $\Omega$ smooth, bounded

- Since K is of order -1 , we know that it is compact, and hence $P:=\frac{1}{2}+\mathrm{K}$ is Fredholm of index $:=\operatorname{dim} T^{-1} 0-\operatorname{dim} T^{*-1} 0=0$.
- Energy estimates ( $\Delta u=0$ and $u \in H_{0}^{1}(\Omega)$ imply $u=0$ ) and the restriction to the boundary (or "jump") relations show that $P$ has zero kernel.
- Hence $P$ is invertible and the solution to our Dirichlet problem is

$$
u(x):=\mathcal{D}\left(P^{-1}(f)\right)=\mathcal{D}\left(\left(\frac{1}{2}+\mathbf{K}\right)^{-1}(f)\right)
$$

Indeed, we already noticed that $u$ is harmonic. Then

$$
\left.u\right|_{\partial \Omega}:=\left.\mathcal{D}\left(\left(\frac{1}{2}+\mathbf{K}\right)^{-1}(f)\right)\right|_{\partial \Omega}=\left(\frac{1}{2}+\mathbf{K}\right)\left(\frac{1}{2}+\mathbf{K}\right)^{-1}(f)=f \mathbb{Q}
$$

## Extensions

This completes our discussion of layer potentials on a smooth, bounded domain $\Omega$.

If $\Omega$ is not smooth or not bounded, K may no longer be compact.
Many further results and extensions: S. Chandler-Wilde, M. Costabel, M. Dauge, M. Kohr, M. Lanza de Cristoforis, V. Mazya, S.E. Mikhailov, M. Putinar, W.L. Wendland, ...

We want to extend these results to the case when $\Omega$ is a domain with cylindrical ends (inside some manifold). (Outlets, wave-guides, tubes (David's talk :-) ... )

Extensions to manifolds: D. \& I. \& M. Mitrea and M. Taylor.

## The Mitrea-Taylor "trick"

Let us assume that $\Omega \subset M=$ a smooth manifold.

The first issue: $\Delta:=d^{*} d$ may no longer be invertible on $M$.
The Mitrea-Taylor "trick" is to replace $\Delta$ with $\Delta+V$, where $V$ is a suitable potential and then, to replace $\Delta^{-1}$ (which may not exist) with $(\Delta+V)^{-1}$ (after proving that the later exists).

Then

$$
\begin{aligned}
& \mathcal{S f}(x):=(\Delta+V)^{-1}\left(f \otimes \delta_{\partial \Omega}\right) \quad \text { and } \\
& \mathcal{D} f(x):=(\Delta+V)^{-1}\left(f \otimes \delta_{\partial \Omega}^{\prime}\right) .
\end{aligned}
$$

## How does the "trick" work?

Let $V \in \mathcal{C}^{\infty}(M)$ satisfy:

- $V \geq 0$ on $M$ and $V=0$ on $\Omega$,
- $V$ is not identically 0 on $M$.

Assume $M$ to be smooth, compact, and connected, then:

- $\Delta+V \geq 0$ is Fredholm of index zero.
- If $(\Delta+V) u=0$, then

$$
0=((\Delta+V) u, u)=(d u, d u)+(V u, u)
$$

and hence $V u=0$ and $d u=0$. Since $V \not \equiv 0, u=0$ ( $L^{2}$-unique continuation property). We obtain that $\Delta+V$ is injective.

- Consequently, $\Delta+V$ is invertible.

The (proved) invertibility of $\Delta+V \geq 0$ allows us to define the layer potential operators

$$
\begin{gathered}
\mathcal{S} f(x):=(\Delta+V)^{-1}\left(f \otimes \delta_{\partial \Omega}\right) \quad \text { and } \\
\mathcal{D} f(x):=(\Delta+V)^{-1}\left(f \otimes \delta_{\partial \Omega}^{\prime}\right)
\end{gathered}
$$

Let us notice that, because $V=0$ inside $\Omega$, we obtain

$$
\Delta \mathcal{D}(f)=(\Delta+V)(\Delta+V)^{-1}\left(f \otimes \delta_{\partial \Omega}^{\prime}\right)=f \otimes \delta_{\partial \Omega}^{\prime}=0
$$

and hence $\mathcal{D}(f)$ (and, similarly, $\mathcal{S}(f)$ ) are again harmonic in $\Omega$, in spite of us having introduced the potential $V$.

The proof that $\frac{1}{2}+\mathbf{K}$ is invertible is then the same.
This then gives the the well-posedness of the Dirichlet problem.

Having explained the Mitrea-Taylor trick for compact manifolds, we now want to extend it to manifolds with cylindrical ends. This will take most of the rest of the first part of the lecture notes.

More precisely,

- We want to first invert $\Delta+V$ (the Mitrea-Taylor trick).
- This allows us to define the layer potential operators. (For instance, recall Df $:=(\Delta+V)^{-1} f$.)
- The jump relations extend immediately (they are the same).
- We then want to show that $\frac{1}{2}+\mathrm{K}$ is also invertible: another use of the Mitrea-Taylor trick.


## The M.-T. trick for cylindrical ends

In order to invert $\Delta+V$ and then $\frac{1}{2}+\mathrm{K}$, we will need:

- To establish that our layer potential operators K and S are "nice" pseudodifferential operators (essentially translation invariant at infinity).
- The Fredholm property for these "nice" pseudodifferential operators.
- The jump relations.
- The $L^{2}$-unique continuation property.

This will be done in the following (several) slides.
(Additional results were given in Mirela Kohr's talk.)

## Outline of this section

## Pseudodifferential operators and cylindrical ends

- In this section, our main interest is in:
- closed manifolds (smooth, compact, without boundary) (0) and
- manifolds with cylindrical ends (1).
- In the second part, more general classes of manifolds, incl.:
- conformally compact manifolds (2);
- Euclidean spaces (3).


## Motivation

- Let $D$ be a differential operator some manifold $M$.
- Assume $D=$ invertible between certain Sobolev spaces.
- its inverse $D^{-1}$ will not be a differential operator. (we assume $D$ is not a multiplication operator).
$D^{-1}=$ a pseudodifferential op. of order $=-\operatorname{order}(D)$ (Beals).
The pseudodifferential operators generalize:
- the differential,
- the convolution, and
- the regularizing operators (order $-\infty$ ).


## The main formula of pseudodifferential theory

The MAIN FORMULA of pseudodifferential theory is:

$$
a(x, D) u(x):=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\imath\langle x, \xi\rangle} a(x, \xi) \hat{u}(\xi) d \xi
$$

(pseudodifferential operator with symbol $a: \mathbb{R}^{n} \times \mathbb{R}^{n} \rightarrow \mathbb{C}$.)
$a(x)$ is a multiplication operator and $a(D)$ is a convolution operator.
Here $\langle x, \xi\rangle:=x_{1} \xi_{1}+x_{2} \xi_{2}+\ldots x_{n} \xi_{n}$ and

$$
\hat{u}(\xi)=\mathcal{F} u(\xi):=\int_{\mathbb{R}} e^{-\imath\langle y, \xi\rangle} u(y) d y .
$$

is the Fourier transform of $u$.
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## Symbols and pseudodifferential operators on $\mathbb{R}^{n}$

Question: Which "symbols" $a(x, \xi)$ should we allow?
Order m, (1,0)-symbols, Hörmander:

$$
\begin{aligned}
& S_{1,0}^{m}\left(\mathbb{R}^{n} \times \mathbb{R}^{n}\right):=\left\{a: \mathbb{R}^{2 n} \rightarrow \mathbb{C} \mid \forall \alpha, \beta, \exists C_{\alpha, \beta} \geq 0\right. \\
&\text { s.t. } \left.\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} a(x, \xi)\right| \leq C_{\alpha, \beta}(1+|\xi|)^{m-|\beta|}\right\} .
\end{aligned}
$$

## Theorem

If $a \in S_{1,0}^{m}\left(\mathbb{R}^{n} \times \mathbb{R}^{n}\right)$, then it defines a continuous map

$$
a(x, D): \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right) \rightarrow \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right) \cdot^{a}
$$

${ }^{a}$ (Schwartz' kernel theorem)

## The third lecture

## The distribution kernel of $a(x, D)$

## Theorem

The distribution kernel $k_{a(x, D)}$ of $a(x, D)$ is:

$$
{ }^{\prime \prime} k_{a(x, D)}(x, y) "=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\imath(x-y, \xi)} a(x, \xi) d \xi=\left(\mathcal{F}_{\xi}^{-1} a\right)(x, x-y) .
$$

Fourier inversion gives $a(x, \xi)$ in terms of $k_{a(x, D)}(x, y)$.
If $\partial \Omega$ is smooth, this proves that S and K are pseudodifferential operators of order $k=-1$ on $\partial \Omega$, because they have distribution kernels

$$
k_{\mathrm{s}}(x, y):=\frac{c_{n}}{|x-y|^{2-n}} \text { and } k_{\mathrm{K}}(x, y):=\frac{c_{n}^{\prime}\left\langle x-y, \nu_{y}\right\rangle}{|x-y|^{2-n}}
$$

which are (essentially) homogeneous of order

$$
-(n-1)-(-1)=-\operatorname{dim}(\partial \Omega)-k .
$$

## Product and diffeomorphism invariance

## Theorem

Let $a \in S_{1,0}^{m}\left(\mathbb{R}^{n} \times \mathbb{R}^{n}\right)$ and $b \in S_{1,0}^{m^{\prime}}\left(\mathbb{R}^{n} \times \mathbb{R}^{n}\right)$. Then
$a b \in S_{1,0}^{m+m^{\prime}}\left(\mathbb{R}^{n} \times \mathbb{R}^{n}\right)$ and there exists $c \in S_{1,0}^{m+m^{\prime}-1}\left(\mathbb{R}^{n} \times \mathbb{R}^{n}\right)$ s. t.

$$
a(x, D) b(x, D)=(a b)(x, D)+c(x, D) .
$$

Localization: Let $U \subset \mathbb{R}^{n}$ be open and

$$
\Psi_{c}^{m}(U):=\left\{a(x, D) \mid \operatorname{supp} k_{a(x, D)} \Subset U \times U\right\} .
$$

## Theorem

For $U \subset \mathbb{R}^{n}$ open, $\Psi_{c}^{\infty}(U):=\cup_{m \in \mathbb{Z}} \Psi_{c}^{m}(U)$ is a filtered algebra invariant under diffeomorphisms.

## Principal symbol

Recall:

- $\Psi_{c}^{m}(U):=\left\{a(x, D) \mid \operatorname{supp}\left(k_{a(x, D)}\right)\right.$ compact $\left.\subset U \times U\right\}$ and
- $a(x, D) b(x, D)=a b(x, D)+$ a lower order operator.

Let $S_{c}^{m}\left(T^{*} U\right):=\left\{a \in S_{1,0}^{m}\left(\mathbb{R}^{n} \times \mathbb{R}^{n}\right) \mid \operatorname{supp}(a) \subset K \times \mathbb{R}^{n}, K \in U\right\}$.

## Theorem

If $a(x, D) \in \Psi_{c}^{m}(U)$, then the principal symbol

$$
\sigma_{m}(a(x, D)):=a+S_{c}^{m-1}\left(T^{*} U\right) \in S_{c}^{m}\left(T^{*} U\right) / S^{m-1}\left(T^{*} U\right)
$$

is well-defined, multiplicative, and diffeomorphism invariant.| UNVERSSIIK |
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## Pseudodifferential operators on manifolds

Let $\phi: U \simeq W \subset \mathbb{R}^{n}, U \subset M$ open (arbitrary $\phi=$ coordinate chart). Diffeomorphism invariance gives

$$
\Psi_{c}^{m}(U):=\phi_{*}^{-1} \Psi_{c}^{m}(W)
$$

## Definition

A linear map $P: \mathcal{C}_{c}^{\infty}(M) \rightarrow \mathcal{C}^{\infty}(M)$ is a pseudodifferential operator of order $\leq m$ on $M$ if, for any $U \subset M$ as above and $\eta \in \mathcal{C}_{c}^{\infty}(U)$, we have

$$
\eta P \eta \in \Psi_{c}^{m}(U)
$$

Let $\Psi^{m}(M)=$ set of order $\leq m$ pseudodifferential operators on $M$.
For instance $\Delta^{-1}=a(D) \in \Psi^{-2}\left(\mathbb{R}^{n}\right)$, although $a(\xi):=-|\xi|^{-1}$ is not in $S_{1,0}^{-2}\left(\mathbb{R}^{n}\right)$ (it has the right decay in $\xi$, but is not smooth).

## Examples and properties

Recall $\Psi^{m}(M)=$ order $m$ pseudodifferential operators on $M$.
If $\partial \Omega$ is smooth, them $\mathbf{S}, \mathrm{K} \in \Psi^{-1}(\partial \Omega)$ since their distribution kernels have the right asymptotic developement at the diagonal and are smooth outside the diagonal.

## Theorem

The principal symbol gives surjective, diffeomorphism invariant maps

$$
\sigma_{m}: \Psi^{m}(M) \rightarrow S_{1,0}^{m} / S_{1,0}^{m-1}\left(T^{*} M\right) .
$$

## Definition

$P \in \Psi^{m}(M)$ is elliptic if $\sigma_{m}(P)$ is invertible (mod. lower order ops.).

## Theorem

Let $a \in S^{m}\left(T^{*} M\right)$. The following are equivalent
(1) There exist $b \in S^{-m}\left(T^{*} M\right)$ such that $a b-1 \in S^{-1}\left(T^{*} M\right)$
(2) There are $C, R>0$ such that $|a(\xi)| \geq C|\xi|^{m}$ for $\xi \in T^{*} M,|\xi| \geq R$.

Then a is called elliptic and this property is a property of its class in $S^{m}\left(T^{*} M\right) / S^{m-1}\left(T^{*} M\right)$.

## Properties of psdos on manifolds

(1) If $P \in \Psi^{m}(M)$, then $P: \mathcal{C}_{c}^{\infty}(M) \rightarrow \mathcal{C}^{\infty}(M)$ continuously, if $P, Q \in \Psi^{\infty}(M)$, the product $P Q$ may not be defined. (It is enough for them to have distribution kernels supported "close" to the diagonal i.e properly supported.)
(2) $\Psi^{m}(M)$ contains all differential operators of order $m$.
(3) $\Psi^{m}(M) \subset \Psi^{m^{\prime}}(M)$ for $m<m^{\prime}$ and $\Psi^{-\infty}(M):=\cap_{m} \Psi^{m}(M)=$ all operators with smooth kernel.
(4) $k_{P}$ is smooth away from the diagonal.

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## Properties of psdos on closed manifolds

Assume $M=$ closed, then we can compose any two pseudos!
We can easily include vector bundles.
(1) If $P \in \Psi^{m}(M ; E)$, then $P: H^{s}(M ; E) \rightarrow H^{s-m}(M ; E)$ is bounded.
(2) If $P \in \Psi^{q}(M ; E), q<0$, then $P: H^{s}(M ; E) \rightarrow H^{s}(M ; E)$ is compact.
(3) If $P \in \psi^{m}(M ; E)$ is elliptic, then, for any $s \in \mathbb{R}$, $P: H^{s}(M ; E) \rightarrow H^{s-m}(M ; E)$ is Fredholm.

The proof of the last result is based on Atkinson's theorem (see next slide).

## Proof of 'elliptic $\Rightarrow$ Fredholm’

## Theorem (Atkinson's theorem)

Let $X$ and $Y$ be Banach spaces; $T \in \mathcal{B}(X ; Y)$ is Fredholm iff there exists $Q \in \mathcal{B}(Y ; X)$ such that $T Q-1_{Y}$ and $Q T-1_{X}$ are compact.

## Proof.

$\mathrm{s}[$ Proof of 'elliptic $\Rightarrow$ Fredholm'] • $P$ is elliptic means that there is $b \in S^{-m}\left(T^{*} M\right)$ such that $\sigma_{m}(P) b-1 \in S^{-1}\left(T^{*} M\right)$.

- The surjectivity of $\sigma_{-m}$ and its multiplicativity give that there exists $Q \in \Psi^{-m}(M)$ (parametrix) such that $P Q-1, Q P-1 \in \Psi^{-1}(M)$.
- Consequently, $P Q-1$ and $Q P-1$ are compact operators. Atkinson's theorem then gives that $P$ is Fredholm.

The converse is true (Fredholm for one simplies elliptic).

## The last two lectures

## "Pictorial" definition of cylindrical ends

We want to extend the Fredholmness result to manifolds with cylindrical ends $M$. (Additional ingredient limit operator.)

- Let $\bar{M}$ be a smooth manifold with boundary $\partial \bar{M}$.
- To $\bar{M}$ we attach the cylindrical end $\partial \bar{M} \times(-\infty, 0]$ along $\partial \bar{M}$.

- Thus $M:=\bar{M} \cup(\partial \bar{M} \times(-\infty, 0])$.


## Translation invariant operators near infinity

Consider the partial translations (isometries)

$$
\begin{gathered}
\Phi_{s}(x, t):=(x, t-s), \quad s \geq 0 \\
\Phi_{s}: \partial \bar{M} \times(-\infty, 0] \rightarrow \partial \bar{M} \times(-\infty,-s] .
\end{gathered}
$$

## Definition

$P: C_{C}^{\infty}(M) \rightarrow C_{C}^{\infty}(M)$ is called translation invariant at infinity if

- its distribution kernel is supported in a neighborhood of the diagonal $\{(x, y) \in M \times M$ : $\operatorname{dist}(x, y)<\varepsilon\}$ for some $\varepsilon>0$, and
- there is $R>0$ s.t. if $\operatorname{supp}(f) \subset \partial \bar{M} \times(-\infty,-R)$ and $s>0$,

$$
P \Phi_{s}(f)=\Phi_{s} P(f) .
$$

Equivalently: $k_{P}\left(t, x, t^{\prime}, x^{\prime}\right)=k_{P}\left(t-s, x, t^{\prime}-s, x^{\prime}\right)$, if $t, t^{\prime} \ll 0$.

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## Psdos translation invariant near infinity

## Definition

Let $\Psi_{\text {inv }}^{m}(M)$ be the space of (classical) pseudodifferential operators $P$ of order $\leq m$ that are translation invariant in a neighborhood of infinity.
(It is contained in the $b$-calculus of Melrose and Schulze.)
As usual:

- $\Psi_{\text {inv }}^{\infty}(M):=\bigcup_{m \in \mathbb{Z}} \Psi_{\text {inv }}^{m}(M)$.
- We can also introduce vector bundles $E, F \rightarrow M$ to obtain

$$
\Psi_{\mathrm{inv}}^{m}(M ; E, F), \ldots
$$

## Almost translation invariant operators at infinity

Drawback: $\Psi_{\text {inv }}^{\infty}(M)$ is not stable under inversion.

## Definition

$\Psi_{\text {ess }}^{\infty}(M)=$ operators essentially translation invariant:

- $\Psi_{\text {ess }}^{m}(M):=\Psi_{\text {inv }}^{m}(M)+\Psi_{\text {ess }}^{-\infty}(M)$
- $\Psi_{\text {ess }}^{-\infty}(M)=$ suitable closure of $\Psi_{\text {inv }}^{-\infty}(M)$ (w.r.t the norms $\|\cdot\|_{m^{\prime}, m}$ of linear operators $H^{m^{\prime}}(M) \rightarrow H^{m}(M), m^{\prime}, m \in 2 \mathbb{Z}$ and with respect to the norms $\left\|a d_{x}^{k}(P)\right\|_{m^{\prime}, m}$.)
$\Psi_{\text {ess }}^{\infty}(M)$ is larger than $\Psi_{\text {inv }}^{\infty}(M)$ and stable under inversion (next). (It is smaller than the c-calculus: Mazzeo-Melrose, Cipriana Anghel, ... )

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## Spectral invariance

- $\Psi_{\text {inv }}^{\infty}(M)$ is a subset of the $\mathbf{b}$-calculus, but $\Psi_{\text {ess }}^{\infty}(M)$ is not.
- Our definition is simple and our algebra is (stable under inversion).


## Theorem (Kohr-Mitrea-V.N.)

Let $T \in \Psi_{\text {ess }}^{m}(M ; E, F) m \geq 0$, be such that $T$ is invertible as a (possibly unbounded ${ }^{a}$ ) operator on $L^{2}(M ; E)$. If $m>0$, we assume also that $T$ is elliptic. Then $T^{-1} \in \psi_{\text {ess }}^{-m}(M ; F, E)$.
${ }^{a} T$ is invertible as an unbounded operator if $T$ is injective and $T^{-1}$ extends to a bounded operator.

## Specific feature: limit operator

## Definition

Let $P \in \Psi_{\text {inv }}^{\infty}(M)$. The limit (or normal) operator associated to $P$ is the unique pseudodifferential operator
$\widetilde{P}: C_{c}^{\infty}(\partial \bar{M} \times \mathbb{R}) \rightarrow C_{c}^{\infty}(\partial \bar{M} \times \mathbb{R})$ s.t.

$$
\widetilde{P}(f):=\Phi_{-s} P \Phi_{s}(f),
$$

for $s$ large enough and $\Phi_{s}$ is the translation by $s$ on $\partial \bar{M} \times \mathbb{R}$.

$$
k_{\tilde{p}}\left(t, x, t^{\prime}, x^{\prime}\right):=\lim _{s \rightarrow \infty} k_{P}\left(t-s, x, t^{\prime}-s, x^{\prime}\right)=k_{P}\left(t-s_{L r g}, x, t^{\prime}-s_{L r g}, x^{\prime}\right) .
$$

- $\widetilde{P}$ is well-defined: $\Phi_{-s} P \Phi_{s}(f)$ is independent of $s$ as above.
- $\widetilde{P}$ is pseudodifferential and translation invariant.


## Fredholm conditions

Next theorem: generalization of the result on closed manifolds.

## Theorem (Kondratiev, Melrose and Mendoza, Kohr-Mitrea-Nistor)

Let $M=\bar{M} \cup(\partial \bar{M} \times(-\infty, 0])$ be a manifold with cylindrical and $P \in \Psi_{\text {ess }}^{m}(M ; E, F)$. Then $P: H^{s}(M ; E) \rightarrow H^{s-m}(M ; F)$ is a Fredholm operator if, and only if,

- it is elliptic and
- its limit operator

$$
\widetilde{P}: H^{s}(\partial \bar{M} \times \mathbb{R} ; E) \rightarrow H^{s-m}(\partial \bar{M} \times \mathbb{R} ; F)
$$

is invertible.
Also, including many generalizations: Anghel, Kondratiev, Mazzeo-Melrose, Lauter-Moroianu, Schrohe, Schulze, Kohr-VN., e. .

## "Pictorial" definition of domains with cylindrical ends

$N$ has a boundary and $M$ does not (both have straight cylindrical ends)


Figure: Manifolds with boundary (or domains) with cylindrical end

The domains with cylindrical ends in $\mathbb{R}^{n}$ are called "with outlets."

## Operators on manifolds with cylindrical ends

- $\Delta=d^{*} d=$ Laplace-Beltrami op. on $M=\bar{M} \cup \partial \bar{M} \times(-\infty, 0]$.
- Then

$$
\Delta=\Delta_{M}=-\partial_{t}^{2}+\Delta_{\partial \bar{M}} \text { on cylindrical end } \partial \bar{M} \times(-\infty, 0],
$$

is translation invariant (in a nbhd of infinity) $\mathbf{v}$ important.

- To construct the layer potential operators $\mathcal{S}, \mathcal{D}, \mathrm{S}$, and K , in the case with boundary, we need to "invert" $\Delta$ (or a suitable modification: Mitrea-Taylor "trick.")
- Recall that inverting $\Delta$ is a substitute for $|x-y|^{2-n}$ (i.e. the convolution with the fundamental solution $c_{n}|x|^{2-n}$ of the Laplacian, which is not possible directly.)


## Doubling and "inverting" $\Delta$

- To construct the layer potential operators on $\Omega=N$ with boundary and cylindrical ends (similar definition, see fig.) we include an additional function $V$ (the Mitrea-Taylor 'trick').

- More precisely, we double $\Omega$ (to obtain $M$ ) by gluing two copies of $\Omega$ along the common boundary and we replace $\Delta$ with $\Delta+V$, where $V$ vanishes on $N$, but not at infinity.
- The limit (jump) properties remain true (this is easy).
- We prove that $\Delta+V$ and then $\frac{1}{2}+K$ are still Fredholm of index zero and injective, hence isomorphisms (using Fredholm and



## Recall how does the "trick" work

Let $V \in \mathcal{C}^{\infty}(M)$ satisfy:

- $V \geq 0$ on $M$ and $V=0$ on $\Omega$,
- $V$ is not identically 0 on $M$ (or at the "infinity" of $M$ ).

Assume $M$ to be smooth, compact, and connected, then:

- $\Delta+V \geq 0$ is Fredholm of index zero.
- If $(\Delta+V) u=0$, then

$$
0=((\Delta+V) u, u)=(d u, d u)+(V u, u)
$$

and hence $V u=0$ and $d u=0$. Since $V \not \equiv 0, u=0$ ( $L^{2}$-unique continuation property). We obtain that $\Delta+V$ is injective.

- Consequently, $\Delta+V$ is invertible and $(\Delta+V)^{-1} \in \Psi_{\text {ess }}^{-2}(M)$

Domains with cylindrical ends: delicate or different issues in red $\mathbb{Q}_{\text {миtuax }}$

## Main result

## Theorem (Mitrea-V.N., Kohr-V.N.)

Let $V \geq 0$ be non-zero and translation invariant at infinity and $\Omega \subset \mathbb{R}^{n}$ be a domain with outlets (cylindrical ends).

- $\Delta+V$ is invertible.
- Let K be the double layer potential operator associated to the $\Delta+V$. Then $\frac{1}{2}+\mathrm{K}$ is invertible and hence the Dirichlet problem on $\Omega$ is well posed.

We need to consider suitable classes of (pseudodifferential) operators that take into account the (translation invariant) structure of the Laplacian on manifolds with cylindrical ends.

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## Proof of the main result

The proof of our main result (the invertibility of $\Delta+V$ and of $\frac{1}{2}+K$ ) is done then as in the compact case (outlined for $\Delta+V$ )

- Fredholm property
- Properties of $\Psi_{\text {ess }}^{m}(M ; E, F)$.
- Jump relations.

Moreover,

$$
(\Delta+V)^{-1} \text { and } \mathrm{K}
$$

are essentially translation invariant (hence also $\left.\left(\frac{1}{2}+K\right)^{-1}\right)$.
Mirela Kohr: generalization to Stokes operator. (Joint w. Wendland.)
Similarities, but also essential differences. (The Neumann problem on domains with conical points behaves quite differently from the Dirichlet problem, Costabel, Dauge, Mazya, Nicaise, V.N., ... )

## Conclusion

Although the double layer potential operator K is no longer compact on a manifold with cylindrical ends, we still obtain the invertibility of $\frac{1}{2}+\mathrm{K}$ and hence the well-posedness of the Dirichlet problem.

We also obtain the structure of the operators $K$ and $\left(\frac{1}{2}+K\right)^{-1}$.

## Thank you for your attention!

The material beginning with next slide were not covered in the Padova course

## A very incomplete bibliography for the first part

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- Petersen, P. Riemannian geometry, 2nd ed., vol. 171 of Grad. Texts Math. New York, NY: Springer, 2006.


# The second part of the lectures notes 

(Not covered in the mini-course.)

## More on Sobolev spaces II

- If $M$ is not compact the Sobolev spaces $W^{m, p}$ depend on the choice of connection on $E$ and on the choice of metric $g$ on $M$, global objects. (Joint with Mirela Kohr.)
- If $M$ is complete, we can define the space $H^{s}(M ; E):=W^{s, 2}(M ; E)$ as the domain of $\left(1-\Delta_{g}\right)^{s / 2}$, where $\Delta_{s}:=-\nabla^{*} \nabla$ (Bochner Laplacian).
- For $s>0$ real, we define $H^{s}(M ; E)$ by interpolation.
- The case $s<0$, for manifold with boundary and relative bounded geometry (later).


## Differential operators

Let $E, F \rightarrow M$ be vector bundles, with $E$ endowed with a connection.

## Definition

Let $a^{[k]}$ be measurable section of $\operatorname{Hom}\left(T^{* \otimes k} \otimes E ; F\right)$ and $\nabla^{0}=i d$. A $\nabla$-differential operator is a map of the form

$$
P=a \cdot \nabla^{\text {tot }}:=\sum_{j=0}^{\mu} a^{[j]} \nabla^{j}: \mathcal{C}^{\infty}(M ; E) \rightarrow \mathcal{C}^{\infty}(M ; F) .
$$

$\operatorname{ord}(P)$ is the least $\mu$ for which such a writing exists (the order of $P$ ).
Suitable extensions by continuity of $P$ will also be called $\nabla$-differential operators and will be denoted by the same letter.

Locally, there is no difference between the $\nabla$-differential operators and the usual differential operators.

## Mapping properties of differential operators

Let $E, F \rightarrow M$ be vector bundles with metrics and metric-compatible with connections.

## Theorem

Let $\ell \in \mathbb{Z}_{+}$and $a^{[k]} \in W^{\ell, \infty}\left(M ; \operatorname{Hom}\left(T^{* \otimes k} \otimes E ; F\right)\right)$ and $\nabla^{0}=i d$. Then the $\nabla$-differential operator

$$
P=a \cdot \nabla^{t o t}:=\sum_{j=0}^{\mu} a^{[j]} \nabla^{j}: W^{s+\mu, p}(M ; E) \rightarrow W^{s, p}(M ; F) .
$$

for $0 \leq s \leq \ell$. (Operator with coefficients in $W^{\ell, \infty}$.)

Better results (including $|s| \leq \ell$ ) for manifolds with bounded geometry, but first the simpler particular case of manifolds with cylindrical ends.

## Curvature

The curvature $R^{E}$ of $(E, \nabla)$ is the anti-symmetric part of $\nabla^{2}$ (i.e. the composite map)

$$
\nabla^{2}: \mathcal{C}^{\infty}(E) \rightarrow \mathcal{C}^{\infty}\left(T^{* \otimes 2} M \otimes E\right) \rightarrow \mathcal{C}^{\infty}\left(\Lambda^{2} T^{*} M \otimes E\right)
$$

That is,

$$
R^{E}(X, Y) u:=\left\langle\nabla^{2} u, X \wedge Y\right\rangle
$$

The curvature of $(M, g)$ is $R=R^{T M}$, the curvature of $E=T M$.

## Curvature

It turns out that the curvature $R^{E}$ is $C^{\infty}(M)$-linear, and hence

$$
R^{E} \in \Omega^{2}(M ; \operatorname{End}(T M))=\mathcal{C}^{\infty}\left(\Lambda^{2} T^{*} M \otimes \operatorname{End}(T M)\right)
$$

Explicitly,

$$
\nabla^{2} u(X, Y)=\nabla_{X} \nabla_{y} u-\nabla_{\nabla_{X} y} u,
$$

and hence the curvature $R^{E}$ of $(E, \nabla)$ is given by

$$
R^{E}(X, Y) u=\nabla_{X} \nabla_{Y} u-\nabla_{Y} \nabla_{X} u-\nabla_{[X, Y]} u .
$$

## Definition

We say that $M$ has totally bounded curvature if its curvatures $R^{M}:=\left(\nabla^{M}\right)^{2}$ and all its covariant derivatives $\left(\nabla^{M}\right)^{k} R^{M}$ are bounded.

## Injectivity radius and geodesics

A $C^{1}$-curve $\gamma:(a, b) \rightarrow M$ is a geodesic if $\nabla_{\gamma^{\prime}(t)}^{M} \gamma^{\prime}(t)=0$. It is locally distance minimizing and uniquely determined by any $\gamma^{\prime}\left(t_{0}\right), t_{0} \in I$. If $M=\mathbb{R}^{d}$ with the usual metric, then a geodesic is just a straight line.

Let

$$
\exp ^{M}(v):=\gamma_{v}(1)
$$

where $\gamma_{v}$ is the unique geodesic with $\gamma^{\prime}(0)=v$, the (geodesic) exponential map.
Let $B_{r}^{X}(x):=\{y \in X \mid d(x, y)<r\}$, and
$r_{\text {inj }}^{M}(p):=\sup \left\{r \mid \exp ^{M}: B_{r}^{T_{p} M}(0) \rightarrow B_{r}^{M}(p)\right.$ is a diffeomorphism $\}$ and

$$
r_{\mathrm{inj}}(M):=\inf _{p \in M} r_{\mathrm{inj}}^{M}(p)
$$

$r_{\text {inj }}(M)=$ the injectivity radius of the Riemannian manifold $(M, g)$.

## Manifolds with bounded geometry

Recall the curvature $R^{M}(X, Y) u=\nabla_{X}^{M} \nabla_{Y}^{M} u-\nabla_{Y}^{M} \nabla_{X}^{M} u-\nabla_{[X, Y]}^{M} u$ and that $M$ has totally bounded curvature if its curvatures $R^{M}:=\left(\nabla^{M}\right)^{2}$ and all its covariant derivatives $\left(\nabla^{M}\right)^{k} R^{M}$ are bounded.
The following concept is crucial.

## Definition

A smooth Riemannian manifold $(M, g)$ is said to have bounded geometry if $M$ has totally bounded curvature and $\operatorname{rinj}(M)>0$.

Then $H^{-s}(M):=H^{-s}(M)^{*}$. (Negative order spaces.)

## Examples

The following are manifolds with bounded geometry:
(1) A closed manifold (i.e. a smooth, compact manifold without boundary).
(1) $M=\mathbb{R}^{d}$ with the standard (Euclidean) metric $d x^{2}:=\left(d x_{1}\right)^{2}+\left(d x_{2}\right)^{2}+\ldots+\left(d x_{d}\right)^{2}$.
(TI) $D \times M$, where $M$ is a manifold with bounded geometry and $D$ is discrete set.
(D) $M_{1} \times M_{2}$, where $M_{1}$ and $M_{2}$ have bounded geometry.
(0) However, if $M$ is a manifold with bounded geometry and $U \subset M$ is an open subset, then $U$ has totally bounded curvature, but may not have positive injectivity radius (so no bounded geometry).

## Boundary and bounded geometry

Let now $M_{0}$ be a Riemannian manifold with boundary, then $r_{\text {inj }}\left(M_{0}\right)=0$, so a manifold with non-empty boundary will never have bounded geometry in the sense of the above definition. The way around this conundrum was found by Schick (2001), who has defined the concept of "manifold with boundary and bounded geometry," (we shall call these manifolds "manifolds with boundary and relative bounded geometry," to avoid confusions). We recall the equivalent definition of manifolds with boundary and relative bounded geometry in (Ammann-Grosse-V.N.)

The main point of that definition is to assume that the boundary $\partial M_{0}$ of $M_{0}$ is a suitable submanifold of a (boundaryless) manifold $M$ with bounded geometry.

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## Second fundamental form

Let hence $M$ be a (boundaryless) manifold with bounded geometry and let us consider a hypersurface $H \subset M$, that is, a submanifold $H$ of $M$ of codimension $\operatorname{dim}(M)-\operatorname{dim}(H)=1$. We assume that $H$ carries a globally defined unit normal vector field $\nu$. We let

$$
\exp ^{\perp}(x, t):=\exp _{x}^{M}\left(t \nu_{x}\right)
$$

be the exponential in the direction of the chosen unit normal vector. We shall need the second fundamental form $\mathrm{II}^{H}$ of $H$ in $M$, which, we recall, is defined by

$$
\Pi^{H}(X, Y) \nu:=\nabla_{X}^{M} Y-\nabla_{X}^{H} Y
$$

where $\nabla^{Z}$ is the Levi-Civita connection of $Z$. Equivalently, since $g\left(\nu, \nabla_{X}^{H} Y\right)=0$, we have $I^{H}(X, Y):=g\left(\nu, \nabla_{X}^{M} Y\right)$.

## Bounded geometry hypersurface

## Definition

Let $(M, g)$ be a Riemannian manifold of bounded geometry and $H \subset M$ be a hypersurface with unit normal vector field $\nu$ on $H$. We say that $H$ is a bounded geometry hypersurface in $M$ if:
(1) $H$ is a closed subset of $M$;
(1) all covariant derivatives $\left(\nabla^{H}\right)^{k} I^{H}, k \geq 0$, are bounded;
(T) $\exp ^{\perp}: H \times(-\delta, \delta) \rightarrow M$ is a diffeomorphism onto its image for some $\delta>0$.

Then $H$ has bounded geometry.
If $H \subset M$ is a compact hypersurface, then $H$ is a bounded geometry hypersurface in $M$.

## Manifolds with boundary and relative bounded geometry

We are ready now to recall the definition of a central concept in analysis on manifolds.

## Definition

We shall say that $M_{0}$ is a manifold with boundary and relative bounded geometry if $M_{0}$ is isometrically contained in a (boundaryless) Riemannian manifold $M$ with bounded geometry such that $\partial M_{0}$ is a bounded geometry hypersurface in $M$.

Note that we use the term "manifold with boundary and relative bounded geometry," which we think is more precise than the term "manifold with boundary and bounded geometry" used before.

## Trace and 'negative' spaces

Let $\partial_{\nu}$ be the normal derivative at the boundary and let

$$
H_{0}^{k}\left(M_{0}\right):=\cap_{j=0}^{k-1} \operatorname{ker}\left(\operatorname{res} \circ \partial_{\nu}^{j}\right)
$$

denote the joint kernel of the restrictions maps res $\circ \partial_{\nu}^{j}, 0 \leq j \leq k-1$.
Theorem (Trace theorem: Grosse-Schneider, Triebel)
Let $M_{0}$ be a manifold with boundary and relative bounded geometry. Then, for every $s>1 / 2$, the restriction res: $\mathcal{C}_{c}^{\infty}\left(M_{0}\right) \rightarrow \mathcal{C}_{c}^{\infty}\left(\partial M_{0}\right)$ extends by continuity to a surjective map

$$
\text { res: } H^{s}\left(M_{0}\right) \rightarrow H^{s-\frac{1}{2}}\left(\partial M_{0}\right)
$$

Moreover, $\mathcal{C}_{c}^{\infty}\left(M_{0} \backslash \partial M_{0}\right)$ is dense in $H_{0}^{k}\left(M_{0}\right)$ and $H^{-k}\left(M_{0}\right)$ identifies with $H_{0}^{k}\left(M_{0}\right)^{*}, k \in \mathbb{N}$.

## Finite width

## Definition

Let $\left(M_{0}, g\right)$ be a Riemannian manifold with boundary $\partial M_{0}$. We say that $M_{0}$ has finite width if:
(1) $\left(M_{0}, g\right)$ is a manifold with boundary and relative bounded geometry and
(1) The function $M_{0} \ni x \rightarrow \operatorname{dist}_{M_{0}}\left(x, \partial M_{0}\right)$ is bounded on $M_{0}$.

The last condition is equivalent to
" $\exists R>0$ such that $M_{0} \subset\left\{x \in M \mid \exists y \in \partial M_{0}, \operatorname{dist}_{M_{0}}(x, y)<R\right\}$."
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## Poincaré inequality

## Theorem (Ammann-Grosse-V.N., Sakurai)

Let $\left(M_{0}, g\right)$ be a Riemannian manifold with finite width. Then there exists $0<C_{M_{0}}<\infty$ such that, for all $f \in \mathcal{C}_{c}^{\infty}\left(M_{0}\right)$ (thus $f=0$ on the boundary of $M_{0}$ ),

$$
\|f\|_{L^{2}\left(M_{0}\right)} \leq C_{M_{0}}\|d f\|_{L^{2}\left(M_{0}\right)} .
$$

A regularity argument then yields the following result.

## Theorem (Ammann-Grosse-V.N.)

Let $M_{0}$ be a smooth Riemannian manifold with smooth boundary $\partial M_{0}$ and finite width. Then $\Delta$ induces isomorphisms

$$
\Delta_{D}=\Delta: H^{m+1}\left(M_{0}\right) \cap H_{0}^{1}\left(M_{0}\right) \rightarrow H^{m-1}\left(M_{0}\right), \quad m \in \mathbb{Z}_{+}:=\{0,1, \ldots\}
$$

## Applications: polygonal regularity

Let us consider a polygonal domain $\Omega_{\infty}$ with maximum angle $\alpha_{M A X}$ and the Poisson problem

$$
\Delta u_{\infty}=f_{\infty} \text { in } \Omega_{\infty}, \quad u_{\infty}=0 \text { on } \partial \Omega_{\infty}
$$

It is known that $u_{\infty}$ has limited regularity, less than $H^{1+\pi / \alpha_{M A X}}$ : Costabel, Dauge, Griesvard, Jerison-Kennig, Kondratiev, Mazya, ...

## Theorem (Limited regularity: Kondratiev, ... )

Let $\delta<\pi / \alpha_{M A X}$. Then there exists $C=C_{\delta}>0$ such that the solution $u_{n} \in H_{0}^{1}\left(\Omega_{n}\right)$ of the equation $\Delta u_{n}=f_{n}$ satisfies

$$
\left\|u_{n}\right\|_{H^{1+\delta}\left(\Omega_{n}\right)} \leq C\left\|f_{n}\right\|_{L^{2}\left(\Omega_{n}\right)} .
$$

## Applications: well-posedness in weighted spaces

The "limited regularity" theorem of the previous slide is not very satisfying. It is, in any case, a consequence of the following more satisfying result in the Babuška-Kondratiev (weighted Sobolev) spaces

$$
\mathcal{K}_{a}^{m}(\Omega):=\left\{\rho^{|\alpha|-a} \partial^{\alpha} u \in L^{2}(\Omega),|\alpha| \leq m\right\}
$$

$\rho=$ distance to vertices (Costabel, Dauge, Kondratiev, Mazya, ... )

## Theorem (Unlimited weighted regularity: Kondratiev '67, ...)

Let $0 \leq \delta<\pi / \alpha_{\text {MAX }}$ and $m \in \mathbb{Z}_{+}$. Then there exists $C=C_{m, \delta}>0$ such that the solution $u_{\infty} \in H_{0}^{1}\left(\Omega_{\infty}\right)$ of the equation $\Delta u_{\infty}=f_{\infty}$ satisfies

$$
\left\|u_{\infty}\right\|_{\mathcal{K}_{\delta+1}^{m+1}\left(\Omega_{\infty}\right)} \leq C\left\|f_{\infty}\right\|_{\mathcal{K}_{\delta-1}^{m-1}\left(\Omega_{\infty}\right)} .
$$

This, as well as the next theorem, are consequences of the well-posedness on manifolds with boundary and relative bounded geometry (Thm 68 stated earlier).

## Applications: "rounding up the corners"

Consider a sequence of smooth domains $\Omega_{n}, n \in \mathbb{Z}$, converging to a polygonal domain $\Omega_{\infty}$ (picture) and the associated Poisson problems $\Delta u_{n}=f_{n}, u_{n} \in H_{0}^{1}\left(\Omega_{n}\right)$, for which we want uniform estimates.


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## Applications: "rounding up the corners"

## Theorem (Daniel-Labrunie-V.N.)

Let $\Omega_{n} \subset \mathbb{R}^{2}$ be a sequence of smooth domains "converging" to a polygonal domain $\Omega_{\infty}$, as in the picture (repeated below). Let $\delta<\pi / \alpha_{M A X}$. Then there exists $C=C_{\delta}>0$ such that for all $n \in \mathbb{N} \cup\{\infty\}$, the solution $u_{n} \in H_{0}^{1}\left(\Omega_{n}\right)$ of $\Delta u_{n}=f_{n}$ satisfies

$$
\left\|u_{n}\right\|_{H^{1+\delta}\left(\Omega_{n}\right)} \leq C\left\|f_{n}\right\|_{L^{2}\left(\Omega_{n}\right)} .
$$


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## The main formula of pseudodifferential theory (repeated)

Recall the MAIN FORMULA of pseudodifferential theory is:

$$
a(x, D) u(x):=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\imath\langle x, \xi\rangle} a(x, \xi) \hat{u}(\xi) d \xi
$$

(pseudodifferential operator with symbol $a: \mathbb{R}^{n} \times \mathbb{R}^{n} \rightarrow \mathbb{C}$.)
Here $\langle x, \xi\rangle:=x_{1} \xi_{1}+x_{2} \xi_{2}+\ldots x_{n} \xi_{n}$ and

$$
\hat{u}(\xi)=\mathcal{F} u(\xi):=\int_{\mathbb{R}} e^{-\imath\langle y, \xi\rangle} u(y) d y .
$$

is the Fourier transform of $u$.

## First example

Let $a: \mathbb{R}^{2 n} \rightarrow \mathbb{C}, a(x, \xi)=a(x)$, (no $\xi$ dependence). Then:

$$
\begin{aligned}
\underline{a(x, D) u(x)} & :=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\imath\langle x, \xi\rangle} a(x) \hat{u}(\xi) d \xi \\
& =a(x) \frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\imath\langle x, \xi\rangle} \hat{u}(\xi) d \xi=\underline{a(x)} u(x),
\end{aligned}
$$

by the Fourier inversion formula:

$$
\left(\mathcal{F}^{-1} w\right)(x):=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}} e^{i\langle x, \xi\rangle} w(\xi) d \xi
$$

## Second example

Let $a: \mathbb{R}^{n} \times \mathbb{R}^{n} \rightarrow \mathbb{C}$ be now $a(x, \xi)=a_{j}(\xi)=\imath \xi_{j}:=\sqrt{-1} \xi_{j}$, with no $x$ dependence ${ }^{1}$ this time. Then:

$$
\begin{aligned}
\begin{array}{l}
a_{j}(D) \\
\end{array}(x)=a(x, D) u(x) & :=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\imath\langle x, \xi\rangle} \xi_{j} \hat{u}(\xi) d \xi \\
& =\frac{\partial u(x)}{\partial x_{j}}=: \underline{\partial_{j}} u(x),
\end{aligned}
$$

because the Fourier transf. interchanges multiplication by $\imath \xi_{j}$ with $\partial_{j}$ :

$$
\mathcal{F} \partial_{j}=\imath \xi_{j} \mathcal{F} \Rightarrow \partial_{j} u=\mathcal{F}^{-1}\left(\imath \xi_{j} \hat{u}\right)=\mathcal{F}^{-1}\left(a_{j} \mathcal{F} u\right) .
$$

## Third example

Iterating, we get

$$
\partial^{\alpha} u:=\partial_{1}^{\alpha_{1}} \partial_{2}^{\alpha_{2}} \ldots \partial_{n}^{\alpha_{n}} u=\mathcal{F}^{-1}\left((\imath \xi)^{\alpha} \hat{u}\right)
$$

Hence, for $a(x, \xi):=\sum_{|\alpha| \leq m} a_{\alpha}(x)(\imath \xi)^{\alpha}$, polynomial in $\xi$ :

$$
\begin{aligned}
a(x, D) u(x) & :=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\imath\langle x, \xi\rangle} \underbrace{\sum_{\alpha \alpha \mid \leq m} a_{\alpha}(x)(\imath \xi)^{\alpha}}_{a(x, \xi)} \hat{u}(\xi) d \xi \\
& =\sum_{|\alpha| \leq m} a_{\alpha}(x) \partial^{\alpha} u(x),
\end{aligned}
$$

a differential op. (We have used $|\alpha|:=\sum \alpha_{j}$.)

## Example $3 \frac{1}{2}$

In the particular case $a(\xi)=1+|\xi|^{2}:=1+\xi_{1}^{2}+\ldots+\xi_{n}^{2}$ :

$$
a(D) u(x):=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\ell\langle x, \xi\rangle}\left(1+|\xi|^{2}\right) \hat{u}(\xi) d \xi=(1-\Delta) u(x) .
$$

It turns out that the inverse of $a(D)$ is of a similar form!

$$
(1-\Delta)^{-1} u(x)=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\imath\langle x, \xi\rangle} \frac{1}{1+|\xi|^{2}} \hat{u}(\xi) d \xi=: b(D) u(x),
$$

where $b(x, \xi)=b(\xi):=\frac{1}{1+|\xi|^{2}}{ }^{2}$
${ }^{2} b(D)=$ convolution with the fundamental solution of $1-\Delta$.

## Symbols and pseudodifferential operators on $\mathbb{R}^{n}$ (repeated)

Recall the order m, (1,0)-symbols, Hörmander:

$$
\begin{aligned}
\boldsymbol{S}^{m}\left(\mathbb{R}^{2 n}\right)=S_{1,0}^{m}\left(\mathbb{R}^{2 n}\right):= & \left\{a: \mathbb{R}^{n} \times \mathbb{R}^{n} \rightarrow \mathbb{C} \mid \forall \alpha, \beta, \exists \boldsymbol{C}_{\alpha, \beta} \geq 0\right. \\
& \text { s.t. } \left.\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} a(x, \xi)\right| \leq C_{\alpha, \beta}(1+|\xi|)^{m-|\beta|}\right\} .
\end{aligned}
$$

Also recall:

## Theorem

If $a \in S^{m}\left(\mathbb{R}^{2 n}\right)$, then it defines a continuous map

$$
a(x, D): \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right) \rightarrow \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right) \cdot \cdot^{a}
$$

[^0]
## Schwartz' kernel theorem

- Let $\langle\rangle:, \mathcal{C}_{c}^{\infty}(M)^{\prime} \times \mathcal{C}_{c}^{\infty}(M) \rightarrow \mathbb{C}$ be the paring between distributions and test functions.
- A linear map $T: \mathcal{C}_{c}^{\infty}(M) \rightarrow \mathcal{C}_{c}^{\infty}(M)^{\prime}$ is continuous if $\langle T \phi, \psi\rangle$ is continuous with respect to $\phi \in \mathcal{C}_{c}^{\infty}(M)$, for any $\psi \in \mathcal{C}_{c}^{\infty}(M)$.


## Theorem (Schwartz' kernel theorem)

Let $T: \mathcal{C}_{c}^{\infty}(M) \rightarrow \mathcal{C}_{c}^{\infty}(M)^{\prime}$ be linear and continuous. Then there exists a unique $k_{T} \in \mathcal{C}_{C}^{\infty}(M \times M)^{\prime}$ such that

$$
\langle T \phi, \psi\rangle=\left\langle k_{T}, \psi \boxtimes \phi\right\rangle
$$

where $(\psi \boxtimes \phi)(x, y)=\psi(x) \phi(y) .\left(k_{T}=\right.$ distribution kernel.)
The converse is also true. Question: Find the kernel $k_{a(x, D)}$.

## Equivalent form using the Fourier transform

If in the main formula

$$
a(x, D) u(x):=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\imath\langle x, \xi\rangle} a(x, \xi) \hat{u}(\xi) d \xi
$$

we substitute the Fourier transform

$$
\hat{u}(\xi)=\mathcal{F} u(\xi):=\int_{\mathbb{R}} e^{-i\langle y, \xi\rangle} u(y) d y,
$$

we obtain the second main formula

$$
a(x, D) u(x)=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}}\left(\int_{\mathbb{R}^{n}} e^{\imath\langle x-y, \xi\rangle} a(x, \xi) u(y) d y\right) d \xi .
$$

(1)

## The distribution kernel of $a(x, D)$

## Theorem

The distribution kernel $k_{a(x, D)}$ of $a(x, D): \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right) \rightarrow \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right) \subset \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)^{\prime}$ is:

$$
" k_{a(x, D)}(x, y) "=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\imath\langle x-y, \xi\rangle} a(x, \xi) d \xi=\left(\mathcal{F}_{\xi}^{-1} a\right)(x, x-y) .
$$

## Proof.

The second main formula gives (integrating wrt $\xi$ ):

$$
\begin{array}{r}
\langle a(x, D) u, v\rangle=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} \underbrace{\int_{\mathbb{R}^{n}} \int_{\mathbb{R}^{n}} e^{\imath\langle x-y, \xi\rangle} a(x, \xi) u(y) d y d \xi}_{a(x, D) u(x)} v(x) d x \\
=\int_{\mathbb{R}^{n}} \int_{\mathbb{R}^{n}} k_{a(x, D)}(x, y) v(x) u(y) d y d x=\left\langle k_{a(x, D)}, v \boxtimes u\right\rangle,
\end{array}
$$

## Properties of the kernel $k_{a(x, D)}$

- The Fourier inversion formula allows us to recover $a(x, \xi)$ from

$$
k_{a(x, D)}(x, y)=\frac{1}{(2 \pi)^{n}} \int_{\mathbb{R}^{n}} e^{\imath\langle x-y, \xi\rangle} a(x, \xi) d \xi
$$

Localization: Let $U \subset \mathbb{R}^{n}$ be open and

$$
\begin{gathered}
\Psi_{c}^{m}(U):=\left\{a(x, D) \mid \operatorname{supp} k_{a(x, D)} \Subset U \times U\right\} \text { and } \\
S_{c}^{m}\left(T^{*} U\right):=\left\{a \in S^{m}\left(\mathbb{R}^{2 n}\right) \mid \operatorname{supp} a \subset K \times \mathbb{R}^{n}, K \Subset U\right\} .
\end{gathered}
$$

The last point shows that

$$
a(x, D) \in \Psi_{c}^{m}(U) \Rightarrow a \in S_{c}^{m}\left(T^{*} U\right) \subset S^{m}\left(\mathbb{R}^{2 n}\right)
$$

## Diffeomorphism invariance

Recall $\Psi_{c}^{m}(U):=\left\{a(x, D) \mid \operatorname{supp} k_{a(x, D)}\right.$ compact $\left.\subset U \times U\right\}$.

## Theorem

For $U \subset \mathbb{R}^{n}$ be open, $\Psi_{c}^{\infty}(U):=\cup_{m \in \mathbb{Z}} \Psi_{c}^{m}(U)$ is an algebra invariant under diffeomorphisms* and under adjoints.
*More precisely, a diffeomorphism $\phi: U \rightarrow W \subset \mathbb{R}^{n}$ induces bijections:

- $\phi^{*}: T^{*} W \rightarrow T^{*} U$, (where $\phi^{*}(x, \xi)=\left(\phi^{-1}(x),(d \phi)^{T} \xi\right)$ );
- $\phi_{*}: \mathcal{C}_{c}^{\infty}(U) \rightarrow \mathcal{C}_{c}^{\infty}(W)$, (where $\phi_{*} u:=u \circ \phi^{-1}$ ); and
- $\phi_{*}: S_{c}^{m}\left(T^{*} U\right) \rightarrow S_{c}^{m}\left(T^{*} W\right)$, (where $\left.\phi_{*} a:=a \circ \phi^{*}\right)$.

Let $a(x, D) \in \Psi_{c}^{m}(U)$, then $\phi_{*} a:=a \circ \phi^{*} \in S_{c}^{m}\left(T^{*} W\right)$ and

$$
\phi_{*} \circ a(x, D) \circ \phi_{*}^{-1}-\left(\phi_{*} a\right)(x, D) \in \Psi_{c}^{m-1}(W) .
$$

## Principal symbol

Let $\phi: U \rightarrow W$ be a diffeomorphism, as before. Then

$$
\begin{gathered}
S_{c}^{m}\left(T^{*} U\right) \ni a \xrightarrow{\phi_{*}} \phi_{*}(a) \in S_{c}^{m}\left(T^{*} W\right) \\
\downarrow \\
\Psi_{c}^{m}(U) \ni a(x, D) \longrightarrow \phi_{*} a(x, D) \phi_{*}^{-1} \in \Psi_{c}^{m}(W)
\end{gathered}
$$

commutes up to lower order symbols.

## Theorem

If $a(x, D) \in \Psi_{c}^{m}(U)$, then the principal symbol

$$
\sigma_{m}(a(x, D)):=a+S^{m-1}\left(T^{*} U\right) \in S^{m}\left(T^{*} U\right) / S^{m-1}\left(T^{*} U\right)
$$

is well-defined, multiplicative, and diffeomorphism invariant.

## Pseudodifferential operators on manifolds (repeated)

Let $M$ be a smooth manifold. Recall the following. If $\phi: U \simeq W \subset \mathbb{R}^{n}$, $U \subset M$ is open, the diffeomorphism invariance allows us to define

$$
\begin{gathered}
\Psi_{c}^{m}(U):=\phi_{*}^{-1} \Psi_{c}^{m}(W) \phi_{*} \quad \text { and } \\
\sigma_{m}: \Psi_{c}^{m}(U) \rightarrow S_{c}^{m} / S_{c}^{m-1}\left(T^{*} U\right)
\end{gathered}
$$

## Definition

A linear map $P: \mathcal{C}_{c}^{\infty}(M) \rightarrow \mathcal{C}^{\infty}(M)$ is a pseudodifferential operator of order $\leq m$ on $M$ if, for any $U \simeq W \subset \mathbb{R}^{n}, U \subset M$, and $\eta \in \mathcal{C}_{C}^{\infty}(U)$, we have $\eta P \eta \in \Psi_{c}^{m}(U)$.

Let $\psi^{m}(M)$ be the set of all pseudodifferential operators on $M$. Then the principal symbol maps extend to surjective maps

$$
\sigma_{m}: \Psi^{m}(M) \rightarrow S^{m} / S^{m-1}\left(T^{*} M\right)
$$

## Pseudodifferential operators on $\mathbb{R}^{n}$

For the particular case $M=\mathbb{R}^{n}$, the space $\Psi^{m}\left(\mathbb{R}^{n}\right)$ contains:
(1) all $a(x, D)$ with $a \in S^{m}\left(\mathbb{R}^{2 n}\right)=S^{m}\left(T^{*} \mathbb{R}^{n}\right)$.
(2) all differential ops. with smooth coefficients.
(3) all operators with smooth distribution kernel.

Not all these operators are of the form $a(x, D)$, with $a \in S^{m}\left(\mathbb{R}^{2 n}\right)$.

## Asymptotic sums

## Theorem

Let $a_{j} \in S^{m_{j}}(M)$, with $m_{j}$ decreasing to $-\infty, j=0,1,2, \ldots$ Then there exists $a \in S^{m_{0}}(M)$ such that, for all $N \in \mathbb{N}, a-\sum_{j=0}^{N} a_{m_{j}} \in S^{m_{N+1}}(M)$.

We then write $a \sim \sum_{j=0}^{\infty} a_{m_{j}}$.

## Definition

We say that $a_{k}: T^{*} M \rightarrow \mathbb{C}$ is (eventually) homogeneous of order $k$ if

$$
a_{k}(t \xi)=t^{k} a_{k}(\xi) \text { for } t,|\xi| \geq 1 .
$$

An eventually homogeneous function of order $m$ is in $S^{m}(M)$ if it is smooth enough.

## Classical symbols

## Definition

A symbol $a \in S^{m}\left(\mathbb{R}^{n}\right)$ is classical if it can be expanded asymptotically $a \sim \sum_{j=0}^{\infty} a_{m-j}$, where $a_{k} \in S^{k}(M)$ is homogeneous of order $k$.

- All results remain true for classical symbols.
- Let $\Psi_{c l}^{m}(M)$ be the corresponding classical psedodifferential op, then $\sigma_{m}: \Psi_{c l}^{m}(M) \rightarrow S_{c l}^{m} / S_{c l}^{m-1}\left(T^{*} M\right)$ has kernel $\Psi_{c l}^{m-1}(M)$.
- If $S^{*} M$ is the set of vectors of length 1 in $T^{*} M$, then

$$
S_{C l}^{m} / S_{C l}^{m-1}\left(T^{*} M\right) \simeq \mathcal{C}^{\infty}\left(S^{*} M\right)
$$

and we can choose $a_{m}$ a representative of $\sigma_{m}(a)$.

- Parametrices of differential operators are classical (index theoty)


## Vector bundles

## We can include vector bundles:

- If $E \rightarrow M$ is a vector bundle, let $E \subset \mathbb{C}^{N}$ be a smooth embedding, and let $e \in \mathcal{C}^{\infty}\left(M ; M_{N}(\mathbb{C})\right)$ be the projection $\mathbb{C}^{N} \rightarrow E$.
- $\Psi^{m}(M ; E):=e M_{N}\left(\Psi^{m}(M)\right) e$ acts on $H^{s}(M ; E) \simeq e H^{s}(M)^{N}$.
- $\sigma_{m}: \Psi^{m}(M ; E) \rightarrow S^{m} / S^{m-1}\left(T^{*} M\right.$; End $\left.(E)\right)$ and the Fredholm theorem remains unchanged.
- Moreover,

$$
S_{c l}^{m} / S_{c l}^{m-1}\left(T^{*} M ; \operatorname{End}(E)\right) \simeq \mathcal{C}^{\infty}\left(S^{*} M ; \operatorname{End}(E)\right),
$$

where $S_{c l}^{m} \subset S^{m}$ denotes the set of classical symbols, as before.

## Outline of the rest of this presentation

We use Lie groupoids, Lie algebroids, and $C^{*}$-algebras to study operators on Lie manifolds generalization of manifolds with cylindrical ends. These objects are motivated by the case of manifolds with cylindrical ends.

This material is about $2 / 3$ of a sequence of lectures (course) delivered in Bruxelles in 2023. The material on pseudodifferential operators is also from that course.

## Motivation: recall manifolds with cylindrical ends

A model class of non-compact manifolds $M$ :

$M$ contains a part (end) that contains a copy of $\partial \bar{M} \times(-\infty, 0]$ with a product metric isometrically: $M=\bar{M} \cup(\partial \bar{M} \times(-\infty, 0])$.

We want now to take a quick look at the differential operators that are invariant (in a neighborhood of) infinity.

## Going back in the definition of with cylindrical ends



Kondratiev's transform $r:=e^{t}$ :

$$
\begin{aligned}
(-\infty, \ln \epsilon) \times \partial \bar{M} \ni(t, y) & \rightarrow(r, y) \in(0, \epsilon) \times \partial \bar{M} \\
\partial_{t} & \rightarrow r \partial_{r}
\end{aligned}
$$

a transl. inv. $\rightarrow$ a smooth on $\bar{M}$.

## Differential operators translation invariant operators near infinity

Back to $\bar{M}$, we get differential operators of the form

$$
P=\sum_{j, Q} a_{Q}\left(e^{t}, y\right)\left(r \partial_{r}\right)^{j} Q(y)
$$

- $Q(y)$ differential operator on the boundary $\partial \bar{M} \ni y$.
- $a_{Q}\left(e^{t}, y\right)=a_{Q}(y)$ for $r=e^{t}$ very small.
- $\widetilde{P}=\sum_{j, Q} a_{Q}(0, y) \partial_{t}^{j} Q(y)$ on $\partial \bar{M} \times \mathbb{R}$.

Diff. ops. enerated by $r \partial_{r}$ and $\partial_{y}$, all tangent to the boundary $\partial \bar{M}$. a Lie algebra.

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## Differential operators in general

- On a closed manifold $M$, we considered all differential operators (all vector fields and all smooth functions).
- In the non-compact case, we need to restrict the choices of vector fields and of smooth functions in order to obtain meaningful results.
- The choices: The 'Quantization Program,' vector fields on manifolds with corners (Cordes, Melrose, Schulze, ... )
- The non-smooth case is treated using the non-compact case, by looking at the set of smooth points $M$ and then choosing a suitable compactification $\bar{M}$ of that set (of smooth points).

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## Manifolds with corners

A manifold with corners $\bar{M}$ is locally of the form $[0,1]^{n}$.

- A face $H \subset \bar{M}$ of maximal dimension is called a hyperface.
- The boundary $\partial \bar{M}$ of $\bar{M}$ is the union of all faces of $\bar{M}$ other than its interior.
- A defining function of a hyperface $H$ of $\bar{M}$ is a function $r$ such that $H=\{r=0\}$ and $d r \neq 0$ on $H$.
- The hyperface $H \subset \bar{M}$ is embedded if it has a defining function.


## Embedded and non-embedded faces

"TEAR-DROP DOMAIN"



$$
\begin{aligned}
& \text { EXACTLY ONE } \\
& \text { HYPERFACE, } \\
& \text { NON EMBEDDED. }
\end{aligned}
$$

## A word on assumptions and notation

- Our vector bundles $E \rightarrow Z$ will always be smooth (and our spaces will be manifolds, possibly with corners).
- $\mathcal{C}^{\infty}(Z ; E)$ denotes the space of smooth sections $s: Z \rightarrow E$ of $E$, also $\mathcal{C}^{\infty}(E)$, if $Z$ is clear from the context.
- $\mathcal{C}_{c}^{\infty}(Z ; E)$ or $\mathcal{C}_{c}^{\infty}(E)$ denotes the space of compactly supported, smooth sections of $E$.
- The Lie bracket will be denoted by [, ].


## The 'Quantization Program'

## Main players:

- $\bar{M}=$ compact manifold with corners.
- $\mathcal{V}_{b}(\bar{M}):=\left\{X \in \mathcal{C}^{\infty}(\bar{M} ; T \bar{M})\right.$ tangent to $\left.\partial \bar{M}\right\}$.
- We will consider a $\mathcal{C}^{\infty}(\bar{M})$ - submodule $\mathcal{V} \subset \mathcal{V}_{b}(\bar{M})$ which is stable for the Lie bracket of vector fields (Lie subalgebra).
- Diff $(\mathcal{V})=$ differential operators generated by $\mathcal{V}$ and $\mathcal{C}^{\infty}(\bar{M})$.


## The quantization program:

To study the analytic properties of the differential operators in Diff $(\mathcal{V})$.
Albin, Mazzeo, Melrose, Piazza, Rochon, (the list is v. incomplete ...)
Convenient approach: using Lie algebroids.

## Lie algebroids (informal definition)

A Lie algebroid is a real vector bundle $A \rightarrow \bar{M}$ (a manifold with corners, as usual) such that " $C^{\infty}(A)$ is a $\mathcal{C}^{\infty}(\bar{M})$-Lie algebra."

More precisely, we are given:
(1) A Lie algebra structure on the space $\mathcal{C}^{\infty}(A)=\mathcal{C}^{\infty}(\bar{M} ; A)$ of global (smooth) sections of $A$.
(2) A Lie algebra action of $\mathcal{C}^{\infty}(A)$ on the base ring $\mathcal{C}^{\infty}(\bar{M})$ :

$$
\mathcal{C}^{\infty}(A) \times \mathcal{C}^{\infty}(\bar{M}) \ni(X, f) \rightarrow X(f) \in \mathcal{C}^{\infty}(\bar{M})
$$

(3) Compatibility conditions (next).

## Lie algebroids (informal discussion)

Saying that $\mathcal{C}^{\infty}(A)$ is a $\mathcal{C}^{\infty}(\bar{M})$-Lie algebra implies the followng compatibility conditions:

- The action $\mathcal{C}^{\infty}(A) \times \mathcal{C}^{\infty}(\bar{M}) \rightarrow \mathcal{C}^{\infty}(\bar{M})$ is a $\mathcal{C}^{\infty}(\bar{M})$-module morphism:

$$
(f X) g=f(X g)
$$

- The bracket [, ] on $\mathcal{C}^{\infty}(A)$ satisfies the Leibnitz rule:

$$
[X, f Y]=f[X, Y]+(X f) Y
$$

- Saying that action on the base ring is a Lie algebra action means that we have a Lie algebra morphism $\mathcal{C}^{\infty}(A) \rightarrow \mathcal{C}^{\infty}(T \bar{M})$ :

$$
[X, Y] f=X(Y f)-Y(X f)
$$

## Lie algebroids

Formalizing " $C^{\infty}(A)$ is a $\mathcal{C}^{\infty}(\bar{M})$-Lie algebra"

## Definition (Pradines 1967, (Book: Mackenzie))

A Lie algebroid is a real vector bundle $A \rightarrow \bar{M}$ together with
(1) a Lie algebra structure [, ] on $\mathcal{C}^{\infty}(A)$ and
(2) a vector bundle map $\varrho: A \rightarrow T \bar{M}$ (anchor, action), such that,
(3) $\forall X, Y \in \mathcal{C}^{\infty}(A), f \in \mathcal{C}^{\infty}(\bar{M}),[X, f Y]=f[X, Y]+(\varrho(X) f) Y$.

- The action $(X, f) \rightarrow X f:=\varrho(X) f$ is automatically a Lie action.
- $\varrho=\varrho_{*}: \mathcal{C}^{\infty}(\bar{M} ; A) \rightarrow \mathcal{C}^{\infty}(\bar{M} ; T \bar{M})$ is the induced map and we shall usually write $X f:=\varrho(X) f$.


## Examples of Lie algebroids

Let us notice first that if $\varrho=0$, then the Leibnitz rule

$$
[X, f Y]=f[X, Y]+(X f) Y=f[X, Y]
$$

simply states that the Lie bracket is bilinear.

## Example 1. A Lie algebra:

- $\bar{M}=p t$ is reduced to a point and $\mathcal{C}^{\infty}(A)=A$ is a Lie algebra.
- the anchor map $\varrho=0$.

Conditions are satisfied: (1) the Lie structure is given; (2) The action of $\mathcal{C}^{\infty}(A)$ on $\mathcal{C}^{\infty}(\bar{M})$ is trivial (zero) because $\rho=0$; (3) Leibnitz is satisfied since the action is trivial $(\varrho=0)$.

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## Examples of Lie algebroids (cont.)

## Example 2. A manifold I (as a topological space):

Let $\bar{M}=$ any manifold with corners and $A$ any vector bundle. We let:
(1) [, ] $=0$ on $\mathcal{C}^{\infty}(A)$.
(2) The action of $\mathcal{C}^{\infty}(A)$ on $\mathcal{C}^{\infty}(\bar{M})$ is trivial $(\rho=0)$.
(3) The Leibnitz rule is satisfied since $\rho=0$ and the bracket is bilinear (even zero).

Importance: Lie algebroids encompass both Lie algebras and smooth manifolds (if $E=\bar{M} \times \mathbb{C}$ ), they "interpolate" between these two classes.

## Exemples of Lie algebroids (cont.)

## Example 3. A manifold II (tangent space):

This time $\bar{M}=M$ closed manifold (smooth, compact, no corners or boundary) and $A=T M$.

- The usual Lie structure (bracket) on $\mathcal{C}^{\infty}(A)=\mathcal{C}^{\infty}(T M)$.
- $\mathcal{C}^{\infty}(T M)$ acts on $\mathcal{C}^{\infty}(M)$ by derivations, that is

$$
\varrho=i d: T \bar{M} \rightarrow T \bar{M} .
$$

- The Leibnitz rule is satisfied by the Lie bracket of vector fields.

This example corresponds to the first example (0) of the last lecture: closed manifolds. The simplest setting in the quantization program is then:
"To study the differential operators on $M$, a closed manifold."

## The Serre-Swan theorem and our main examples

The next example correspond to the second example discussed in the last lecture: manifolds with cylindrical ends (1).

To make the connection(s):

## Theorem (Serre-Swan)

Let $\mathcal{E}$ be a projective, finitely generated $\mathcal{C}^{\infty}(\bar{M})$-module. Then there exists a smooth, finite-dimensional vector bundle $E \rightarrow \bar{M}$ such that $\mathcal{E} \simeq \mathcal{C}^{\infty}(\bar{M} ; E)$ naturally as $\mathcal{C}^{\infty}(\bar{M})$-modules.

The converse is also true (but not used in this talk).
( $\bar{M}$ is a manifold with corners, as usual).

## The b-calculus and cylindrical ends

## Example 4. The $b$-calculus (or cylindrical ends)

- Let $\bar{M}=$ a manifold with boundary $\partial \bar{M}$.
- $\mathcal{V}_{b}(\bar{M}):=\left\{X \in \mathcal{C}^{\infty}(\bar{M} ; T \bar{M})|X|_{\partial \bar{M}}\right.$ tangent to $\left.\partial \bar{M}\right\}$.
- Near the boundary of $\bar{M}$, we can assume $\bar{M}=[0, \infty) \times \mathbb{R}^{n-1}$. Then $\mathcal{V}_{b}(\bar{M})$ is even a free module with basis

$$
x_{1} \partial_{1}, \partial_{2}, \partial_{3}, \ldots, \partial_{n}
$$

- Serre-Swan theorem gives that $\mathcal{V}_{b}(\bar{M})$ identifies with the sections of a vector bundle usually denoted ${ }^{b} T \bar{M} .\left(\mathcal{V}=\mathcal{V}_{b}(\bar{M})\right)$.
- ${ }^{b} T \bar{M}$ is a Lie algebroid because $\mathcal{V}_{b}(\bar{M})$ is a Lie algebra.

The Laplacian in generalized spherical coordinates corresponds to the this class of Lie algebroids. (Conical points, APS, ... )

## The "edge"-calculus

## Example 5. Mazzeo's "edge"-calculus:

- Let $\bar{M}=$ a manifold with boundary plus a fibration $\partial \bar{M} \rightarrow B^{k}$
- $\mathcal{V}_{e}:=\left\{X|X|_{\partial \bar{M}}\right.$ tangent to the fibers of $\left.\partial \bar{M} \rightarrow B^{k}\right\}$.
- Near the boundary of $\bar{M}$, we can assume $\bar{M}=[0, \infty) \times \mathbb{R}^{n-1}$. Then $\mathcal{V}_{e}(M)$ is even a free module with basis

$$
x_{1} \partial_{1}, \partial_{2}, \ldots, \partial_{n-k}, x_{1} \partial_{n-k+1}, \ldots, x_{1} \partial_{n}
$$

- Serre-Swan theorem gives that $\mathcal{V}_{e}(\bar{M})$ identifies with the sections of a vector bundle ${ }^{e} T \bar{M}$ (Mazzeo $\mathcal{V}=\mathcal{V}_{e}(\bar{M})$ ).
- ${ }^{e} T \bar{M}$ is a Lie algebroid because $\mathcal{V}_{e}$ is a Lie algebra.
- $\partial \bar{M} \rightarrow B^{k}:=p t: b$-calculus $(k=0)$.
- $\partial \bar{M} \rightarrow B^{k}:=\partial \bar{M}$ : conformally compact manifolds $(k=n-1)$.
- The Laplacian in cylindrical coordinates: $n=3, k=1$.


## Lie algebroids and Geometry

Many examples of Lie algebroids often arise from Geometry:

- Foliations.
- This time we have a subbundle $A=\mathcal{F} \subset T M$ such that $\mathcal{C}^{\infty}(M ; \mathcal{F})$ is stable for the Lie bracket (and hence a Lie algebra).
- In particular, $\mathcal{C}^{\infty}(A)=\mathcal{C}^{\infty}(\mathcal{F})$ consists of vector fields tangent to the leaves of the foliation.
- Again, $\varrho: \mathcal{F} \rightarrow$ TM is the identity (more precisely, the inclusion).
(Connes. Also Androulidakis, Benameur, Skandalis, ... ).
- Poisson manifolds. $A=T^{*} M \rightarrow M$.

ATTENTION: the vector associated to some $X \in \mathcal{C}^{\infty}(A)$ may vanish at some point, without $X$ vanishing at that point. Ex: Poisson not symplectic. ISOTROPY (next).

## Isotropy $\mathfrak{g}_{x}$

Let $\varrho: A \rightarrow T \bar{M}$ be a Lie algebroid over $\bar{M}$ and $x \in \bar{M}$. We set:

- $\mathfrak{g}_{x}:=\operatorname{ker}\left(\varrho_{x}: A_{x} \rightarrow T_{x} \bar{M}\right)$; (finite dim.)
- $\mathcal{V}:=\mathcal{C}^{\infty}(A)$, a Lie algebra;
- $\mathcal{V}_{x}:=\left\{Y \in \mathcal{V} \mid \varrho_{x}(Y)=0\right\}=\left\{Y \in \mathcal{V} \mid Y(x) \in \mathfrak{g}_{x}\right\}$.
- $I_{x}:=\left\{f \in \mathcal{C}^{\infty}(\bar{M}) \mid f(x)=0\right\}$; and

If $Y \in \mathcal{V}_{x}$ and $f \in \mathcal{C}^{\infty}(\bar{M})$, then $(Y f)(x)=0$, as well.
Let $Y \in \mathcal{V}$, then

$$
[X, f Y]=X(f) Y+f[X, Y] \in I_{X} \mathcal{V}
$$

and therefore $I_{x} \mathcal{V}$ is an ideal in $\mathcal{V}_{x}$.
Consequently, $\mathfrak{g}_{x}=\mathcal{V}_{x} / I_{x} \mathcal{V}$ is a Lie algebra. It is the isotropy of $A$ at $x$ and will play a very, very important role in what follows.

## Lie algebroids and isotropy (quick review)

Summary: A Lie algebroid (Pradines 1967) is a real vector bundle $A \rightarrow \bar{M} \mathrm{~s}$. t.

$$
" \mathcal{C}^{\infty}(A) \text { is a } \mathcal{C}^{\infty}(\bar{M}) \text {-Lie algebra." }
$$

More precisely:
(1) $\mathcal{V}:=\mathcal{C}^{\infty}(A)$ is a Lie algebra.
(2) $\mathcal{V}:=\mathcal{C}^{\infty}(A)$ acts on $\mathcal{C}^{\infty}(\bar{M})$.
(3) Compatibility conditions (Leibnitz rule, ... ).

We usually give the action: $\varrho: A \rightarrow T \bar{M}$ (anchor map).

## Isotropy $\mathfrak{g}_{x}$ of $\varrho: A \rightarrow T \bar{M}$

## Theorem

Let $x \in \bar{M}$, then $\mathfrak{g}_{x}:=\operatorname{ker}\left(\varrho_{x}: A_{x} \rightarrow T_{x} \bar{M}\right)$ is a Lie algebra.

## Proof.

$\mathfrak{g}_{x}=\mathcal{V}_{x} / I_{x} \mathcal{V}$, where

- $\mathcal{V}_{x}:=\left\{Y \in \mathcal{C}^{\infty}(A) \mid \varrho_{x}(Y)=0\right\}$.
- $I_{x}:=\left\{f \in \mathcal{C}^{\infty}(\bar{M}) \mid f(x)=0\right\}$.

Leibnitz' rule gives that $I_{x} \mathcal{V}$ is an ideal in $\mathcal{V}_{x}$.
Consequently, $\mathfrak{g}_{x}=\mathcal{V}_{x} / I_{x} \mathcal{V}$ is a Lie algebra.
$\mathfrak{g}_{x}$ is the isotropy of $A$ at $x$ and will play a very, very important role.

## Examples of isotropies

## Example 1. A Lie algebra:

- $A=$ Lie algebra, $\bar{M}=p t, \varrho=0$.
- The isotropy is $\mathfrak{g}_{x}=A$, maximal.


## Example 2. A manifold I (as a topological space):

(1) $0=\varrho: A \rightarrow T \bar{M}$ a vector bundle, $[]=$,0 on $\mathcal{C}^{\infty}(A)$.
(2) The isotropy is $\mathfrak{g}_{x}=A_{x}$, maximal at each point.

## Example 3. A manifold II (tangent space):

- $A=T M \rightarrow T M$ with the Lie bracket and $\varrho=i d$.
- The isotropy is $\mathfrak{g}_{x}=0$, minimal at each point.


## The isotropies of the $b$-calculus

## Example 4. $b$-calculus (or cylindrical ends)

- $\bar{M}=$ a manifold with boundary and $A={ }^{b} T \bar{M}$, so that $\mathcal{C}^{\infty}(A)=\mathcal{V}_{b}(\bar{M}):=\left\{X \in \mathcal{C}^{\infty}(T \bar{M})|X|_{\partial \bar{M}}\right.$ tangent to $\left.\partial \bar{M}\right\}$.
- Near the boundary, $\mathcal{V}_{b}(\bar{M})$ is a free module with basis

$$
x_{1} \partial_{1}, \partial_{2}, \partial_{3}, \ldots, \partial_{n}
$$

- Let $x \in M:=\bar{M} \backslash \partial \bar{M}$, then $\rho_{x}=i d$, so $\mathfrak{g}_{x}=0$, minimal.
- Let $x \in \partial \bar{M}$, then $\mathfrak{g}_{x}=\mathbb{R}\left(x_{1} \partial_{1}\right)$, NO LONGER minimal.
- The section $x_{1} \partial_{1}$ of $A={ }^{b} T \bar{M}$ vanishes nowhere, but its associated vector field $\varrho\left(x_{1} \partial_{1}\right)=x_{1} \partial_{1}$ vanishes at the boundary. Notice that $\partial_{1}$ does not define a section of $A={ }^{b} T \bar{M}$.


## The isotropies of the "edge"-calculus

## Example 5. Mazzeo's "edge"-calculus:

- $\bar{M}=$ a manifold with boundary, $\partial \bar{M} \rightarrow B^{k}$ smooth fibration.
- $\mathcal{C}^{\infty}(A)=\mathcal{V}_{e}:=\left\{X|X|_{\partial \bar{M}}\right.$ tangent to the fibers of $\left.\partial \bar{M} \rightarrow B^{k}\right\}$.
- Near the boundary, $\mathcal{V}_{e}(M)$ is a free module with basis

$$
x_{1} \partial_{1}, \partial_{2}, \ldots, \partial_{n-k}, x_{1} \partial_{n-k+1}, \ldots, x_{1} \partial_{n} .
$$

- In the interior the isotropy vanishes, but if $x \in \partial \bar{M}$, then the isotropy $\mathfrak{g}_{x}$ is linearly generated by the sections

$$
x_{1} \partial_{1}, x_{1} \partial_{n-k+1}, \ldots, x_{1} \partial_{n} .
$$

- $\mathfrak{g}_{x}$ is a non-commutative Lie algebra if $k>0$. (The semi-direct product $\mathbb{R}^{k} \rtimes \mathbb{R}$.) More difficult!


## Outline

## We now introduce a generalization of manifolds with cylindrical ends following Ammann-Lauter-V.N.

## Definition of 'Lie manifolds'

Recall $\mathcal{V}_{b}(\bar{M}) \subset \mathcal{C}^{\infty}(\bar{M} ; T \bar{M})=$ vector fields tangent to all faces.

## Definition (Ammann-Lauter-V.N.)

A Lie manifold is pair $(\bar{M}, \mathcal{V})$ consisting of a compact manifold with corners $\bar{M}$ and a subspace $\mathcal{V} \subset \mathcal{V}_{b}(\bar{M})$ of vector fields s.t.
(1) $\mathcal{V}$ is closed under the Lie bracket [, ];
(2) $\mathcal{V}$ is a finitely-generated, projective $\mathcal{C}^{\infty}(\bar{M})$-module;
(3) $\mathcal{C}_{c}^{\infty}(M ; T M) \subset \mathcal{V}($ recall $M:=\bar{M} \backslash \partial \bar{M})$.
(1) and (2) $\Leftrightarrow$ Lie algebroid $A \rightarrow \bar{M}$. Addl. prop. (3).

Examples 3-5 (compact, b-calc. and "edge"-calc.).

## Compatible Riemannian metric

Let $(\bar{M}, \mathcal{V})$ be a Lie manifold.

- $\mathcal{V}=$ projective $\mathcal{C}^{\infty}(\bar{M})$-module implies that there exists a vector bundle $A$ s.t. $\mathcal{V} \simeq \mathcal{C}^{\infty}(\bar{M} ; A)$. (The Serre-Swan theorem.)
- $A=$ Lie algebroid (because $\mathcal{V}=$ Lie algebra.) $(M, A)$ also called a Lie manifold.
- The inclusion $\mathcal{C}_{c}^{\infty}(M ; T M) \subset \mathcal{V}$ is equivalent to $A=T M$ in the interior $M:=\bar{M} \backslash \partial \bar{M}$ of $\bar{M}$.
- Any metric on $A$ will induce a metric on $T M$ (i.e. on $M$ ), called compatible metric, and is unique up to Lipschitz equivalence. It is complete and has positive injectivity radius.


## Properties of compatible metrics

Recall $T M \subset A$, where $M:=\bar{M} \backslash \partial \bar{M}$. The Levi-Civita connection

$$
\nabla: \mathcal{C}^{\infty}(T M) \rightarrow \mathcal{C}^{\infty}\left(T^{*} M \otimes T M\right)
$$

associated to a compatible metric $g$ on $M$ (comming from a metric on $A$ ) extends to a map ( $A$-connexion, Weinstein and col.)

$$
\nabla: \mathcal{C}^{\infty}(A) \rightarrow \mathcal{C}^{\infty}\left(A^{*} \otimes A\right) .
$$

(Proof: Koszul's formula ... )
Theorem (Ammann-Lauter-V.N.)
All geometric differential operators associated to a compatible metric $g$ are in $\operatorname{Diff}(A)$.

For instance, the Laplacian $\Delta_{g} \in \operatorname{Diff}(A)$.
(Also for the other operators: we need vector bundles.)

## Bounded geometry +

- The curvature $\nabla^{2} \in \mathcal{C}^{\infty}\left(\bar{M} ; A^{* \otimes 2} \times \operatorname{End}(A)\right)$ is bounded.
- Bounded geometry (also positive injectivity radius, non-trivial).
- If $E, F \rightarrow \bar{M}$ are vector bundles, their "right connections" are the $A$-connections $\nabla: \mathcal{C}^{\infty}(\bar{M} ; A) \rightarrow \mathcal{C}^{\infty}\left(\bar{M} ; A^{*} \otimes A\right)$.
- Differential operators $\operatorname{Diff}(A ; E, F)$ generated by $\nabla$ and $\operatorname{Hom}(E ; F)$.
- Clifford bundles $W \rightarrow \bar{M}$, with a metric and $A$-connection, $c: \boldsymbol{A}^{*} \rightarrow \operatorname{End}(W)$, with $c(\xi) c(\eta)+c(\eta) c(\xi)=2\langle\xi, \eta\rangle$ I, then the associated Dirac operator $Q^{W} \in \operatorname{Diff}(A ; W)$ (Parker's talk).


## Sobolev spaces

- $(M, A)=$ a Lie manifold $\left(\mathcal{V}=\mathcal{C}^{\infty}(A)\right)$.
- $g=$ a compatible metric on $M:=\bar{M} \backslash \partial \bar{M}$ (from a metric on $A$ ).
- $L^{p}(M)$ is defined wrt the measure induced by $g$.
- If $s \in \mathbb{R}_{+}$, we let $H^{s}(M)$ be the domain of $\left(1-\Delta_{g}\right)^{s / 2}$.
- If $m \in \mathbb{N}$,

$$
H^{m}(M):=\left\{u \mid X_{1} X_{2} \ldots X_{k} u \in L^{2}(M), k \leq m, X_{j} \in \mathcal{V}\right\}
$$

Kondratiev 1967, Mazya, Mazzeo, Melrose, Plamenevskij, ... (Ammann-Lauter-V.N.)

## First example (simplest): cylindrical ends

- $\bar{M}=$ a manifold with smooth boundary $\partial \bar{M}=\{x=0\}$.
- $\mathcal{V}=\mathcal{V}_{b}(\bar{M})=$ vector fields on $\bar{M}$ tangent to $\partial \bar{M}$ (recall projective).
- No cond. on vector fields in the interior (all Lie man).
- Choice of compatible metric (near the boundary):

$$
\frac{\left(d x_{1}\right)^{2}}{x_{1}^{2}}+h\left(x^{\prime}\right), \quad h \text { metric on } \partial \bar{M} \ni x^{\prime}
$$

Kondratiev's transform $x=e^{t}$ transforms this metric into

$$
(d t)^{2}+h\left(x^{\prime}\right)
$$

the Riemannian metric of a manifold with cylindrical ends.
Pseudodifferential calculus: $b$-calculus.
it contains the 'inv'-calculus of the first lecture. (Isotropy $\mathfrak{g}_{x} \simeq \mathbb{R}$.)

## Second example: asymptotically hyperbolic manifolds

- As before, $\bar{M}$ with smooth boundary $\partial \bar{M}=\{x=0\}$.
- $\mathcal{V}=\mathcal{V}_{0}=x \mathcal{C}^{\infty}(\bar{M} ; T \bar{M})=$ the space of vector fields on $\bar{M}$ that vanish on the boundary.
- Particular case of the "edge"-calculus $(\partial \bar{M} \rightarrow \partial \bar{M})$.
- Local basis: $x \partial_{x}, x \partial_{x_{2}}, \ldots, x \partial_{x_{n}}$ (so projective module).
- Compatible metric: $g=\frac{h}{x^{2}}$, where $h=$ metric on $\bar{M}$.

Metric: conformally compact (ex: asymptotically hyperbolic).
Pseudodifferential calculus: Lauter, Mazzeo, Schulze.
(Isotropy $\mathfrak{g}_{x} \simeq \mathbb{R}^{n-1} \rtimes \mathbb{R}$ at the boundary.)

## Third example: asymptotically Euclidean manifolds

- As before, $\bar{M}$ with smooth boundary $\partial \bar{M}=\{x=0\}$.
- $\mathcal{V}=x \mathcal{V}_{b}(\bar{M})=$ the space of vector fields on $\bar{M}$ that vanish on the boundary $\partial \bar{M}$ and whose normal covariant derivative to the boundary also vanishes.
- Local basis $x_{1}^{2} \partial_{x}, x \partial_{x_{2}}, \ldots, x \partial_{x_{n}}$.
- $\mathfrak{g}_{x}=\mathbb{R}^{n}$ at the boundary (trivial inside, all Lie Man.).

The resulting metric and geometry is that of an asymptotically conical manifold (includes the asymptotically Euclidean case).
If $\bar{M}=$ spherical compactification of $\mathbb{R}^{n}$, the resulting Riemannian manifold is simply $\mathbb{R}^{n}$ with the usual metric. ('SG' or sc calculus, Melrose, Parenti, Schrohe, ... ) $N$-body problems (Georgescu, Vasy, Ammann-Mougel-V.N.)

## Picture of the examples

$$
r^{\alpha} \partial_{r}:
$$

$$
t= \pm \int \frac{d r}{r \alpha}
$$

## "No number" example

The metric on $M=$ interior of $\bar{M}$ does not determine $\bar{M}$.
Example, $M=\mathbb{R}^{n}$ with the Euclidean metric:

- $\bar{M}_{1}:=\overline{\mathbb{R}}^{n}$, radial compactification, as in the previous example.
- $\bar{M}_{2}:=\overline{\mathbb{R}}^{k} \overline{\mathbb{R}}^{n-k}$, product of radial compactifications.
(Also $\bar{M}_{3}:=(\overline{\mathbb{R}})^{n}$, a product of "closed" lines.)
- $N$-body problems (Kottke, Vasy, Ammann-Mougel-V.N.).

Different differential operators, different analytic properties.

## Comparison of the three examples of Lie manifolds

$h=$ (non-singular, true) metric on $\bar{M}$. Examples of compatible metrics:
(1) Cylindrical ends: $\mathcal{V}=\mathcal{C}^{\infty}(\bar{M}) r \partial_{r}+\sum \mathcal{C}^{\infty}(\bar{M}) \partial_{y}$,

$$
g=\frac{(d r)^{2}}{r^{2}}+h
$$

(2) Conformally compact: $\mathcal{V}=\mathcal{C}^{\infty}(\bar{M}) r \partial_{r}+\sum \mathcal{C}^{\infty}(\bar{M}) r \partial_{y}$,

$$
g=\frac{(d r)^{2}+h}{r^{2}}=\frac{h_{0}}{r^{2}}
$$

(3) Asymptotically conical: $\mathcal{V}=\mathcal{C}^{\infty}(\bar{M}) r^{2} \partial_{r}+\sum \mathcal{C}^{\infty}(\bar{M}) r \partial_{y}$,

$$
g=\frac{(d r)^{2}}{r^{4}}+\frac{h}{r^{2}}
$$

## Picture of the examples

$$
d t:= \pm \frac{d r}{\pi} ; \quad \text { a. cyl. \& hyp.: } \alpha=1 ; \quad \text { a. euclidean } \alpha=2 \text {. }
$$



## Examples: $\Delta$ in polar and cylindrical coordinates

(1) $r^{2} \Delta_{\mathbb{R}^{2}}=\left(r \partial_{r}\right)^{2}+\partial_{\theta}^{2}$ is the differential operator generated by

$$
r \partial_{r} \quad \text { and } \quad \partial_{\theta}
$$

on $\bar{M}=[0, \infty) \times S^{1}$. (Totally characteristic operators.)
(2) Similarly, $r^{-2} \Delta_{\mathbb{R}^{3}}=\left(r \partial_{r}\right)^{2}+\partial_{\theta}^{2}+\left(r \partial_{z}\right)^{2}$ is the differential operator generated by

$$
r \partial_{r}, \quad \partial_{\theta}, \quad \text { and } \quad r \partial_{z}
$$

on $\bar{M}:=[0, \infty) \times S^{1} \times \mathbb{R}$. Edge differential operators.

## Limit operators for the $b$-calculus

$b$-calculus:
( $\bar{M}=$ compact manifold with boundary.
$\mathcal{V}_{b}(\bar{M}):=\left\{X \in \mathcal{C}^{\infty}(\bar{M} ; T \bar{M})\right.$ tangent to $\left.\partial \bar{M}\right\}$ and $\operatorname{Diff}\left(\mathcal{V}_{b}(\bar{M})\right)$ of differential operators generated by $\mathcal{V}_{b}(\bar{M}) . r=x$ distance to the boundary $\partial \bar{M}$.)

Let $x_{1}=r, x_{2}, \ldots, x_{n}$ be local coordinates near some boundary point.
Then, $P \in \operatorname{Diff}\left(\mathcal{V}_{b}(\bar{M})\right)$ if, and only if,

$$
P u(x)=\sum_{|\alpha| \leq m} a_{\alpha}(x)\left(x_{1} \partial_{1}\right)^{\alpha_{1}} \partial_{2}^{\alpha_{2}} \ldots \partial_{n}^{\alpha_{n}} u(x)
$$

(Totally characteristic operators. Never elliptic if $\partial \bar{M} \neq \emptyset$.) Indicial operator $r \partial_{r} \rightarrow \partial_{t}$ related to the isotropy!
(1) UNIVERSITÉ UNIVERSIE

## Fredholm conditions

- $(\bar{M}, \mathcal{V})=$ Lie manifold and let $\{Z\}$ be the orbits of $\mathcal{V}$ on $\partial \bar{M}$.
- We associate to each orbit $Z$ a simply-connected Lie group $G_{z}$, with Lie $G_{z} \simeq \mathfrak{g}_{x}, x \in Z$, and
- associate to each $D \in \operatorname{Diff}(\mathcal{V})$ a differential operator $D_{z}$ a $G_{Z}$-invariant differential operator on $Z \times G_{Z}$ (limit operator).


## Conjecture (Carvalho-V.N.-Qiao, ... )

Let $(\bar{M}, \mathcal{V})$ be a Lie manifold and $D \in \operatorname{Diff}(\mathcal{V})$ have order $m$.
$D: H^{s}(M) \rightarrow H^{s-m}(M)$ is Fredholm $\Leftrightarrow$
$D$ elliptic and all $D_{Z}: H^{s}\left(Z \times G_{z}\right) \rightarrow H^{s-m}\left(Z \times G_{z}\right)$ are invertible .

## Back to the quantization program

Say $D \in \operatorname{Diff}(A)$ is of order $2 m$ and we want to prove that

$$
D: H^{2 m}(M) \rightarrow L^{2}(M)
$$

is invertible (Hadamard well-posed) or Fredholm.
$\left(1+\Delta_{g}\right)^{m}: H^{2 m}(M) \rightarrow L^{2}(M)$ is an isomorphism (the definition).
Our problem is then equivalent to proving that

$$
D\left(1+\Delta_{g}\right)^{-m}: L^{2}(M) \rightarrow L^{2}(M)
$$

is invertible (resp. Fredholm).
$D\left(1+\Delta_{g}\right)^{-m}$ is a pseudodifferential operator of order zero (no longer differential). (Technical reasons: norm closure: $C^{*}$-algebra; we lose asymptotics, regularity, ... , but OK for Fredholm.)

## Pseudodifferential operators

(Pseudodifferential operators were discussed only briefly.)

- Let $(\bar{M}, A)$ be a Lie manifold, $M:=\bar{M} \backslash \partial \bar{M}$, the interior of $\bar{M}$.


## The 'Quantization Program'

To construct a pseudodifferential calculus $\Psi_{p r}^{m}(A)$ on $M$ with the usual symbolic properties and which reflects the properties of $\operatorname{Diff}(A)$, the differential operators generated by $\mathcal{V}=\mathcal{C}^{\infty}(A)$.

Usually many choices! We will use the fact that $A$ extends $T M$ and $A^{*}$ extends $T^{*} M$.

- We fix a compatible metric $g$ on $M$, as before $(T M \subset A)$.
- Let exp : $T M \rightarrow M$ be the exponential map (associated to $g$ ), which we know exists since $M$ is complete.



## Definition of $\Psi_{p r}^{-\infty}(A)$

Let $0<r<r_{M}$, where $r_{M}$ is the injectivity radius of $M$, (positive).
Let $\pi: A \rightarrow \bar{M}$ be the projection. For $k \in \mathcal{C}^{\infty}\left(M^{2}\right)$ and $v \in T M$ we set

$$
\tilde{k}(v):=k(\pi(v), \exp (v))
$$

Thus $\tilde{k}=k \circ(p, \exp ) \in \mathcal{C}^{\infty}(T M)$.
We let $\Psi_{r}^{-\infty}(A)$ be the operators with kernels $k: M^{2} \rightarrow \mathbb{C}$ with support in $\left\{\operatorname{dist}_{g}(x, y)<r\right\}$ such that $\tilde{k}$ extends to a smooth function on $A$.

## Definition

Finally, we let $\Psi_{p r}^{-\infty}(A)=$ the algebra generated by $\Psi_{r}^{-\infty}(A)$.
(It consists of properly supported ops; unlike the $b$ and edge calculian

## The construction of $\Psi_{p r}^{m}(A)$

To define the "very small calculus" $\Psi_{p r}^{m}(A)$, we choose first a "quantization" map $q: S_{C l}^{\infty}\left(T^{*} M\right) \rightarrow \Psi^{\infty}(M)$ as follows.
Let $\chi:[0, \infty) \rightarrow[0,1]$ be smooth cut-off f. (1 near 0, 0 far from 0 ) and

$$
\begin{gathered}
q: S_{c l}^{m}\left(T^{*} M\right) \rightarrow \psi^{m}(M), \\
{[q(a) u](x):=(2 \pi)^{-n} \int_{T_{x}^{*} M} a(\xi)\left(\int_{T_{x} M} e^{-\imath\langle y, \xi\rangle} \chi(\|y\|) u(\exp (y)) d y\right) d \xi .}
\end{gathered}
$$

## Definition

We then let $\Psi_{p r}^{m}(A):=q\left(S_{c l}^{m}\left(A^{*}\right)\right)+\Psi_{p r}^{-\infty}(A)$.
( $S_{c l}^{m}(E)$ was defined using local trivializations of the v . bundle $E \rightarrow \bar{M}$.)

## Properties

- $\Psi_{p r}^{m}(A) \Psi_{p r}^{m^{\prime}}(A) \subset \Psi_{p r}^{m+m^{\prime}}(A)$.
- $\sigma_{m}(q(a))=a+S_{c l}^{m-1}\left(A^{*}\right)$ for all $a \in S_{c l}^{m}\left(A^{*}\right)$; hence, $D$ is a differential operator in $q\left(S_{c l}^{m}(A)\right)$, iff $D \in \operatorname{Diff}(A)$.
- $D\left(1+\Delta_{g}\right)^{-m} \in \bar{\Psi}^{0}(A)$ in general, but $D\left(1+\Delta_{g}\right)^{-m} \in \bar{\Psi}^{0}(A)=\mathrm{a}$ groupoid $C^{*}$-algebra in favorable situations.
- compact operators $=\mathcal{K} \subset \bar{\Psi}^{0}(A)$ and we obtain a map

$$
\bar{\Psi}^{0}(A) / \mathcal{K} \rightarrow \mathcal{B}\left(L^{2}(M)\right) / \mathcal{K}
$$

which will decide the Fredholm property of our operators. We just need to decide which operators are invertible in $\bar{\psi}^{0}(A) / \mathcal{K}$.

## The role of groupoids (informal)

- Ehresman (Brussels 1959).
- Examples of groupoids $\mathcal{G}$ : Melrose's " $b$-double space"
- The groupoid $\mathcal{G}$ must satisfy $A(\mathcal{G})=A$, its Lie algebroid be the one appearing in the quantization program. (Integration of Lie algebroids!)
- Depending on what one wants, the groupoid must satisfy some additional conditions ("nice")
- If the groupoid is "nice" ( $\mathcal{G} \mid \partial \bar{M}$ Hausdorff and satisfies the (ISF) condition ("Strong Exel property")) $\Rightarrow$ Fredholm conditions.


## Groupoids

## Definition

A groupoid $\mathcal{G}$ is a small category s.t. all morphisms are invertible.

## Notations:

(1) $\mathcal{G} \rightrightarrows \bar{M}$ denotes a groupoid with objects (or units) $M$.
(2) $d, r: \mathcal{G} \rightarrow \bar{M}$ give the domain and range of a morphism.

Typically, $\bar{M}$ is compact with corners and $M$ is its interior.
A more concrete definition of a groupoid $d, r: \mathcal{G} \rightrightarrows \bar{M}$ is in terms of the structural morphisms $d, r, \mu, u, \iota$ that define it.

## Canonical morphisms

The structural morphisms $d, r, \mu, u, \iota$ satisfy:
(1) The prod $\mu(g, h):=g h$ is defined if, and only if, $d(g)=r(h)$;
(2) $\mu: \mathcal{G}^{(2)}:=\{(g, h) \mid d(g)=r(h)\} \rightarrow \mathcal{G}$, associative.
(3) "Unit map" $u: \bar{M} \rightarrow \mathcal{G}$, s.t. $g u(d(g))=g=u(r(g)) g$.
(4) $d(u(x))=r(u(x))=x$ (so we will identify $x \simeq u(x)$.)

Above: $\mathcal{G}=$ category.Next: "all morphisms are invertible."
(5) "Inverse" $\iota: \mathcal{G} \rightarrow \mathcal{G}$, s.t.

$$
g \iota(g)=r(g)=u(r(g)) \text { and } \iota(g) g=d(g)
$$

## Outline

## We show how to use $C^{*}$-algebras, Lie groupoids, their Lie algebroids and associated algebras to characterize Fredholm operators

## Fredholm conditions

- $(\bar{M}, \mathcal{V})=$ Lie manifold and $\{Z\}=$ the orbits of $\mathcal{V}$ on $\partial \bar{M}$.
- Orbit $Z \rightarrow G_{z}$, a simply-connected Lie group with Lie $G_{z} \simeq \mathfrak{g}_{x}$,
- $D \in \operatorname{Diff}(\mathcal{V} ; E, F) \rightarrow$ differential operators $D_{Z}=\pi_{Z}(D)$ a $G_{Z}$-invariant differential op. on $Z \times G_{Z}$ (limit operator).


## Conjecture (Carvalho-V.N.-Qiao, ... )

Let $(\bar{M}, \mathcal{V})$ be a Lie manifold and $D \in \operatorname{Diff}(\mathcal{V} ; E, F)$ have order $m$.
$D: H^{s}(M ; E) \rightarrow H^{s-m}(M ; F)$ is Fredholm $\Leftrightarrow D$ elliptic and all $D_{Z}: H^{s}\left(Z \times G_{Z} ; E\right) \rightarrow H^{s-m}\left(Z \times G_{Z}, F\right)$ are invertible .
$\sigma_{m}(D) \in \mathcal{C}^{\infty}\left(S^{*} A ; \operatorname{Hom}(E, F)\right)$.
Statement about the $\left\{\pi_{z}\right\}, C^{*}$-algebras.

## Proving the Fredholm property

$\mathcal{H}=$ Hilbert space, $\mathcal{K}(\mathcal{H})=$ compact operators on $\mathcal{H}$.

- How to prove that $T \in \mathcal{B}(\mathcal{H})$ is Fredholm?
- Answer: Atkinson's theorem ( $T$ is invertible modulo $\mathcal{K}(\mathcal{H})$ ).
- Equivalent formulation, the image of $T$ is invertible in $\mathcal{B}(\mathcal{H}) / \mathcal{K}(\mathcal{H})$.
- It is convenient to use $C^{*}$-algebras.


## $C^{*}$-algebras

- A concrete $C^{*}$-algebras is a subalgebra $A=A^{*}=\bar{A} \subset \mathcal{B}(\mathcal{H})$.
- If $1, a \in A$ and $a^{-1} \in \mathcal{B}(\mathcal{H})$, then $a^{-1} \in A$.
- An abstract $C^{*}$-algebras is one isometrically isomorphic to a concrete one. (Ex. $\mathcal{B}(\mathcal{H}) / \mathcal{K}(\mathcal{H})$.)
- How to prove that something is invertible in $A$ ?
- "Invertibility sufficient families of representations" (Rabinovich-Roch-Silbermann).


## Invertibility sufficient families of representations

$A=C^{*}$-algebra.

- A representation of $A$ is a *-morphism $\pi: A \rightarrow \mathcal{B}\left(\mathcal{H}_{\pi}\right)$.


## Definition

Let $\mathcal{F}:=\{\pi\}$ be a family of representations of $A$. We say that $\mathcal{F}$ is "invertibility sufficient" (ISF) if the following condition is satisfied:
$a \in A$ is invertible in $A$ if, and only if, $\pi(a)$ is invertible in $\mathcal{B}\left(\mathcal{H}_{\pi}\right)$ for all $\pi \in \mathcal{F}$.

- Example: $A=\mathcal{C}(X)$, where $X$ is a compact space (model commutative $C^{*}$-algebra). We let

$$
\mathcal{F}:=\left\{e_{x}\right\}, \quad x \in X
$$

$e_{x}(f)=f(x)$, the evaluation at $x$. Then $\mathcal{F}$ is (ISF).

## $C^{*}$-algebras

In our applications:

- The family $\{\pi\}$ will be the family of regular representations of some Lie groupoid $\mathcal{G} \mid \partial \bar{M}$ ( where $\mathcal{G}$ is the "double space").
- $\pi(a)$ will be the "limit operators" of $a$.
- The algebra $A=C^{*}(\mathcal{G}) / \mathcal{K}$ or $A=\bar{\Psi}^{0}(\mathcal{G}) / \mathcal{K}$.
- We obtain Fredholm conditions when the family of regular representations of $\mathcal{G} \mid \partial \bar{M}$ is (ISF) condition ("Strong Exel property").


## Lie groupoids and convolution

## Definition

A Lie groupoid $\mathcal{G} \rightrightarrows \bar{M}$ is a groupoid such that $\mathcal{G}=$ manifold (poss. non Hausdorff!), $\bar{M}$ is a Hausdorff manifold (poss. w. corners), all structural maps are smooth, and $d$ is a submersion.

The space $\mathcal{C}_{C}(\mathcal{G})$ becomes an involutive (pre)Banach algebra:

$$
\left(\varphi_{1} * \varphi_{2}\right)(g):=\int_{d(h)=d(g)} \varphi_{1}\left(g h^{-1}\right) \varphi_{2}(h) \mathrm{d} \lambda_{d(g)}(h)
$$

$\lambda_{x}$ is a suitable set of measures on $\mathcal{G}_{x}:=d^{-1}(x)$ (Haar system, it is the natural volume form for Lie groupoids: $\mathcal{G}_{x}=$ manifold).
(norm $\|f\|:=\max _{x \in \bar{M}} \int_{\mathcal{G}_{x}}|f(h)| d \lambda_{x}(h)$. )

## The Lie algebroid of a Lie groupoid

- Let $\mathcal{G} \rightrightarrows \bar{M}$ be a Lie groupoid with structural map $d, r: \mathcal{G} \rightrightarrows \bar{M}$.
- We consider $\mathcal{V}$ to be the space of $d$-vertical vector fields (tangent to the fibers of $d$ ) that are right invariant.
- $\mathcal{V}$ dentifies with the sections of a vector bundle on $\bar{M}$, the Lie algebroid $A(\mathcal{G})$ of $\mathcal{G}$.
- The measures on $\mathcal{G}_{x}:=d^{-1}(x)$ comes from a metric on $A(\mathcal{G})$.

To use $\mathcal{G}$, we need $A(\mathcal{G})=A$, that is, we need to integrate the given Lie algebroid (Lie's third theorem). This is the famous construction of the double spaces (Mazzeo, Melrose, ... , the triple space is the space of composable arrows).
There are general theorems, but the resulting groupoid is useless.©

## Groupoid $C^{*}$-algebras

Define the regular representation $\pi_{x}, x \in \bar{M}$,

$$
\pi_{x}: \mathcal{C}_{c}(\mathcal{G}) \rightarrow \mathcal{B}\left(L^{2}\left(\mathcal{G}_{x}, \lambda_{x}\right)\right), \quad \pi_{x}(f) g:=f * g .
$$

## Definition (Reduced $C^{*}$-algebra $C_{r}^{*}(\mathcal{G})$ )

$C_{r}^{*}(\mathcal{G})$ is the completion of $\mathcal{C}_{C}(\mathcal{G})$ with respect to

$$
\|\varphi\|_{r}:=\sup \left\|\pi_{x}(\varphi)\right\|, \quad x \in \bar{M} .
$$

If $\mathcal{G}$ is metrically amenable (i.e. $C^{*}(\mathcal{G}) \simeq C_{r}^{*}(\mathcal{G})$ ), then we have the short exact sequence for all $U=F^{c}$ open, invariant (Renault)

$$
0 \rightarrow C_{r}^{*}\left(\mathcal{G}_{U}\right) \rightarrow C_{r}^{*}(\mathcal{G})^{+} \rightarrow C_{r}^{*}\left(\mathcal{G}_{F}\right)^{+} \rightarrow 0
$$

We obtain Fredholm conditions when $\mathcal{G} U:=d^{-1}(U)=U \times U$, since $C^{*}\left(\mathcal{G}_{U}\right) \simeq \mathcal{K}$ (next).

## Examples of Lie groupoids I

(1) The pair groupoid $\mathcal{H}:=M \times M$ is the groupoid $\mathcal{H}$ having exactly one arrow between any two units. Hence

$$
d(x, y)=y, r(x, y)=x, \quad \text { and }(x, y)(y, z)=(x, z)
$$

$A(M \times M)=T M$.
The Haar measures $\lambda_{x}$ on $\mathcal{G}_{x}=M \times\{x\}, x \in M$, satisfy $\lambda_{x}=\lambda_{y}=\lambda$

$$
\phi_{1} * \phi_{2}(x, z):=\int_{M} \phi_{1}(x, y) \phi_{2}(y, z) d \lambda(y)
$$

and hence

$$
C^{*}(\mathcal{H}) \simeq C_{r}^{*}(\mathcal{H}) \simeq \mathcal{K}\left(L^{2}(M, \lambda)\right)
$$

Finally, $\pi_{x}$ is the natural representation on $L^{2}(M ; \lambda)$ (indep. of $x$ ).

## Other examples of Lie groupoids II

(3) A locally compact group: $G$ with its Haar measure, $M=\{e\}$.
$C^{*}(G) \cong C_{r}^{*}(G) \Leftrightarrow G$ is amenable (this is the case when we have edge calculi or desingularization in general).
(4) A space: $\mathcal{G}=\bar{M}$ (only units), $C^{*}(\mathcal{G})=C_{r}^{*}(\mathcal{G})=\mathcal{C}_{0}(\bar{M})$.
(5) The product of the two: $\mathcal{G}=\bar{M} \times G \rightrightarrows \bar{M}, d=r$.
(6) A bundle of Lie groups $\mathcal{G} \rightarrow \bar{M}$, in which case $d=r$.

The last example is the main building block for the groupoids in applications.

All these examples satisfy (ISF) if, and only if, the groups involved (if any) are amenable.

## Other examples of Lie groupoids III

(3) The cross-product groupoid: obtained from the action of a locally compact group $G$ on manifold with corners $\bar{M}$.
Then $\mathcal{G}=G \times \bar{M} \rightrightarrows \bar{M}$ with

$$
d(g, x)=x, r(g, x)=g x,(h, g x)(g, x)=(h g, x),
$$

and $C^{*}(\mathcal{G})=\mathcal{C}(\bar{M}) \rtimes G$ and $C_{r}^{*}(\mathcal{G})=\mathcal{C}_{0}(\bar{M}) \rtimes_{r} \mathcal{G}$.
(cross-prod: Exel, Georgescu-Iftimovici, Karlovich, Mantoiu, Mougel, Rabinovich, Roch, Silbermann, ... ).
Used for the "scattering calculus" and for the $N$-body problem with $G=\mathbb{R}^{n}$ and $\bar{M}$ a suitable blow-up of $\overline{\mathbb{R}}^{n}$, (Georgescu, Vasy, Ammann-Mougel-V.N.)

## Other examples of groupoids IV

(8) The "b-groupoid" $\mathcal{G}_{b}(\bar{M})$ for $\bar{M}$ with smooth boundary $\partial \bar{M}$ :

$$
\mathcal{G}_{b}(M):=(M \times M) \sqcup(\partial \bar{M} \times \partial \bar{M}) \times \mathbb{R},
$$

$M:=\bar{M} \backslash \partial \bar{M}$ is the interior of $\bar{M}$.
$\mathcal{G}_{b}(\bar{M})$ is obtained from the 'double-space' construction by removing outside faces

Also Grieser, Lesch, Monthubert, Schrohe, Schulze, ...
$M \subset \bar{M}$ and $\partial \bar{M} \subset \bar{M}$ are invariant subsets of $\bar{M}$. The restriction of $\mathcal{G}_{b}(M)$ to $M$ is $M \times M$ and to $\partial \bar{M}$ is $\partial \bar{M} \times \partial \bar{M} \times \mathbb{R}$. Each satisfies (ISF), so $\mathcal{G}_{b}(M)$ satisfies (ISF), next.

## The (ISF) property

Assume $\mathcal{G}$ to be Hausdorff (non trivial!).
If $U \subset \bar{M}$ is an invariant open subset and if $\left.\mathcal{G}\right|_{U}$ and $\left.\mathcal{G}\right|_{U}$ satisfy (ISF,) then $\mathcal{G}$ satisfies (ISF), a consequence of the exact sequence

$$
0 \rightarrow C_{r}^{*}\left(\mathcal{G}_{U}\right) \rightarrow C_{r}^{*}(\mathcal{G})^{+} \rightarrow C_{r}^{*}\left(\mathcal{G}_{F}\right)^{+} \rightarrow 0 .
$$

We can "slice and dice" our groupoid (which we could not have done with our operator!)

We prove the Fredholm Conjecture in stages (for sufficiently nice groupoids).

We can "glue" groupoids that satisfy the conjecture (Rémy Côme).
If a groupoid $\mathcal{G}$ satisfies the conjecture, then the groupoid associated to a desingularization also satisfies the conjecture ("blow-up invariant"

## The desingularization groupoid I

- Let $\mathcal{G} \rightrightarrows M$ be a Lie groupoid and $L \subset M$ a "nice" submanif, meaning that it has a tubular nbhd in an algebroid sense.
- Let $\pi: S N L \rightarrow L$ be the unit sphere bundle in the normal bundle $\pi: N L \rightarrow L$ of $L$ in $M$ ("pies" everywhwere!).
- Let $[M: L]=$ usual blow-up obtained by replacing $L$ with SNL.
- On $L$ we consider the bundle of Lie groups $\mathcal{L}:=T L \rtimes \mathbb{R}_{+}^{*} \xrightarrow{\pi} L$ and define the fibered pull-back groupoid

$$
\mathcal{H}:=S N L \times_{L} S N L \times_{L} \mathcal{L}:=\{(x, y, v) \mid \pi(x)=\pi(y)=\pi(v)\}
$$

The desingularization is then

$$
\begin{gathered}
{[[\mathcal{G}: L]]:=\mathcal{H} \sqcup \mathcal{G}_{M \backslash L}^{M \backslash L} \rightrightarrows[M: L]} \\
\left(\mathcal{G}_{A}^{B}=d^{-1}(A) \cap r^{-1}(B)\right) . \text { Uses res. of Debord-Skandalis. }
\end{gathered}
$$

## The desingularization groupoid II

## The desingularization

$$
[[\mathcal{G}: L]]:=\mathcal{H} \sqcup \mathcal{G}_{M \backslash L}^{M \backslash L} \rightrightarrows[M: L]
$$

was constructed to satisfy the usual condition ( $r_{L}=$ dist to $L$ ):

$$
\mathcal{C}^{\infty}(A([[\mathcal{G}: L]]))=r_{L} \mathcal{C}^{\infty}([M: L]) \mathcal{C}^{\infty}(A((\mathcal{G}))
$$

If $M$ has no boundary and $\mathcal{G}=M \times M$, then any smooth $L \subset M$ is nice $(A(\mathcal{G})$-tame $)$ and the desingularization $[[\mathcal{G}: L]]$ defines the "edge" calculus for the boundary fibration SNL := $\partial[M: L] \rightarrow L$ (Grušin, Schulze, Mazzeo, Krainer).

If, furthermore, $L$ is a point, we obtain the $b$-groupoid for $[M: L]$.

## Conclusion

- The Fredholm conjecture may be true for all Lie manifolds. Hard (close to conjectures in $C^{*}$-algebras: Effros-Hahn, ...)
- In practice, however: the hardest part is to integrate the Lie algebroid $A$ defining the Lie manifold $\mathcal{C}^{\infty}(A)=\mathcal{V}$, the vector fields. You want a Hausdorff groupoid as "double space."
- There exist general results (Crainic-Fernandez, Debord, V.N.,) but they do not give (in general) Hausdorff groupoids.
- If one is lucky, $\mathcal{G}=\cup_{z} Z \times Z \times G_{Z} \cup M \times M$, (use [V.N.] $\approx 2000$ ).
- Check that on has a stratification of the units $\bar{M} \supset \mathcal{O}_{k}$, with the orbits $Z$ closed in $\mathcal{O}_{k+1} \backslash \mathcal{O}_{k}$ and same stabilizer. (Ex. Compact group actions: Albin-Melrose.) True in all examples that I know.


## Thank you for your attention!

See the last slide for some references.

## References for the second part

This list is very incomplete, but may help you get started.
[1] Pseudodifferential operators on differential groupoids (with Weinstein and Xu ) https://arxiv.org/pdf/funct-an/9702004.pdf
[2] Groupoids and the integration of Lie algebroids https://arxiv.org/pdf/math/0004084.pdf
[3] Fredholm conditions on non-compact manifolds: theory and examples (with Carvalho and Qiao) https://arxiv.org/pdf/1703.07953.pdf
[4] Gluing groupoids: Fredholm conditions and layer potentials (Carvalho, Côme, and and Qiao) https://arxiv.org/pdf/1811.07699.pdf

## REMOVED STUFF 1

## "Pictorial" definition of cylindrical ends (repeated)

A model class of non-compact manifolds: the manifolds with (straight) cylindrical ends (picture):


As for manifolds with bounded geometry, the Sobolev spaces on manifolds with cylindrical ends can be defined using partitions of unity.

The formal definition is on the next slides.

## Formal definition of $M$ with cylindrical ends



- $\bar{M}=$ a smooth, compact Riemannian manifold, $\partial \bar{M} \neq \emptyset$.
- To $\bar{M}$ we attach the cyl. end $\partial \bar{M} \times(-\infty, 0]$ along $\partial \bar{M} \equiv \partial \bar{M} \times\{0\}$ :

$$
M:=\bar{M} \cup \partial \bar{M} \times(-\infty, 0]
$$

## Translation invariant operators near infinity

Recall the partial translations (isometries) $\Phi_{s}(x, t):=(x, t-s)$, $s \geq 0$ on the half-infinite cylindrical end

$$
\Phi_{s}: \partial \bar{M} \times(-\infty, 0] \rightarrow \partial \bar{M} \times(-\infty,-s] .
$$

## Definition

$P: C_{C}^{\infty}(M) \rightarrow C_{C}^{\infty}(M)$ is called translation invariant at infinity if

- its distribution kernel is supported in a neighborhood of the diagonal $\{(x, y) \in M \times M: \operatorname{dist}(x, y)<\varepsilon\}$ for some $\varepsilon>0$, and
- there is $R>0$ s.t. if $\operatorname{supp}(f) \subset \partial \bar{M} \times(-\infty,-R)$ and $s>0$,

$$
P \phi_{s}(f)=\phi_{s} P(f) .
$$

The differential operators that are translation invariant at infinity have coefficients in $W^{\infty, \infty}$, and hence, they are bounded on all Sobolev spaces. The same is true of the pseudodifferential operators.

## Differential operators

Let $E, F \rightarrow M$ be vector bundles, with $E$ endowed with a connection.

## Definition

Let $a^{[k]}$ be measurable section of $\operatorname{Hom}\left(T^{* \otimes k} \otimes E ; F\right)$ and $\nabla^{0}=i d$. A $\nabla$-differential operator is a map of the form

$$
P=a \cdot \nabla^{\text {tot }}:=\sum_{j=0}^{\mu} a^{[j]} \nabla^{j}: \mathcal{C}^{\infty}(M ; E) \rightarrow \mathcal{C}^{\infty}(M ; F) .
$$

$\operatorname{ord}(P)$ is the least $\mu$ for which such a writing exists (the order of $P$ ).
Suitable extensions by continuity of $P$ will also be called $\nabla$-differential operators and will be denoted by the same letter.

Locally, there is no difference between the $\nabla$-differential operators and the usual differential operators.

## Mapping properties of differential operators

Let $E, F \rightarrow M$ be vector bundles with metrics and metric-compatible with connections.

## Theorem

Let $\ell \in \mathbb{Z}_{+}$and $a^{[k]} \in W^{\ell, \infty}\left(M ; \operatorname{Hom}\left(T^{* \otimes k} \otimes E ; F\right)\right)$ and $\nabla^{0}=i d$. Then the $\nabla$-differential operator

$$
P=a \cdot \nabla^{t o t}:=\sum_{j=0}^{\mu} a^{[j]} \nabla^{j}: W^{s+\mu, p}(M ; E) \rightarrow W^{s, p}(M ; F) .
$$

for $0 \leq s \leq \ell$. (Operator with coefficients in $W^{\ell, \infty}$.)

Better results (including $|s| \leq \ell$ ) for manifolds with bounded geometry, but first the simpler particular case of manifolds with cylindrical ends.

## Curvature

It turns out that the anti-symmetric part of $\nabla^{2}$ (i.e. the composite map)

$$
\nabla^{2}: \mathcal{C}^{\infty}(E) \rightarrow \mathcal{C}^{\infty}\left(T^{* \otimes 2} M \otimes E\right) \rightarrow \mathcal{C}^{\infty}\left(\Lambda^{2} T^{*} M \otimes E\right)
$$

is actually linear (curvature, see also next).

## Curvature

## The curvature

$$
R \in \Omega^{2}(M ; \operatorname{End}(T M))=\mathcal{C}^{\infty}\left(\Lambda^{2} T^{*} M \otimes \operatorname{End}(T M)\right)
$$

of $(E, \nabla)$ (or of $(M, g)$ if $E=T M$ ) is such that

$$
\left\langle\nabla^{2} u, X \wedge Y\right\rangle=R(X, Y) u
$$

Explicitly, $\nabla^{2} u(X, Y)=\nabla_{X} \nabla_{Y} u-\nabla_{\nabla_{X} Y} u$, and hence the curvature $R$ of $(E, \nabla)$ is given by

$$
R(X, Y) u=\nabla_{X} \nabla_{Y} u-\nabla_{Y} \nabla_{X} u-\nabla_{[X, Y]} u .
$$

## Definition

We say that $M$ has totally bounded curvature if its curvatures $R^{M}:=\left(\nabla^{M}\right)^{2}$ and all its covariant derivatives $\left(\nabla^{M}\right)^{k} R^{M}$ are bounded.

## Injectivity radius and geodesics

A $C^{1}$-curve $\gamma:(a, b) \rightarrow M$ is a geodesic if $\nabla_{\gamma^{\prime}(t)}^{M} \gamma^{\prime}(t)=0$. It is locally distance minimizing and uniquely determined by any $\gamma^{\prime}\left(t_{0}\right), t_{0} \in I$. If $M=\mathbb{R}^{d}$ with the usual metric, then a geodesic is just a straight line.

Let

$$
\exp ^{M}(v):=\gamma_{v}(1)
$$

where $\gamma_{v}$ is the unique geodesic with $\gamma^{\prime}(0)=v$, the (geodesic) exponential map.
Let $B_{r}^{X}(x):=\{y \in X \mid d(x, y)<r\}$, and
$r_{\text {inj }}^{M}(p):=\sup \left\{r \mid \exp ^{M}: B_{r}^{T_{p} M}(0) \rightarrow B_{r}^{M}(p)\right.$ is a diffeomorphism $\}$ and

$$
r_{\mathrm{inj}}(M):=\inf _{p \in M} r_{\mathrm{inj}}^{M}(p)
$$

$r_{\text {inj }}(M)=$ the injectivity radius of the Riemannian manifold $(M, g)$.

## Manifolds with bounded geometry

Recall the curvature $R^{M}(X, Y) u=\nabla_{X}^{M} \nabla_{Y}^{M} u-\nabla_{Y}^{M} \nabla_{X}^{M} u-\nabla_{[X, Y]}^{M} u$ and that $M$ has totally bounded curvature if its curvatures $R^{M}:=\left(\nabla^{M}\right)^{2}$ and all its covariant derivatives $\left(\nabla^{M}\right)^{k} R^{M}$ are bounded.
The following concept is crucial.

## Definition

A smooth Riemannian manifold $(M, g)$ is said to have bounded geometry if $M$ has totally bounded curvature and $\operatorname{rinj}(M)>0$.

Then $H^{-s}(M):=H^{-s}(M)^{*}$. (Negative order spaces.)

## Examples

The following are manifolds with bounded geometry:
(1) A closed manifold (i.e. a smooth, compact manifold without boundary).
(1) $M=\mathbb{R}^{d}$ with the standard (Euclidean) metric $d x^{2}:=\left(d x_{1}\right)^{2}+\left(d x_{2}\right)^{2}+\ldots+\left(d x_{d}\right)^{2}$.
(TI) $D \times M$, where $M$ is a manifold with bounded geometry and $D$ is discrete set.
(D) $M_{1} \times M_{2}$, where $M_{1}$ and $M_{2}$ have bounded geometry.
(0) However, if $M$ is a manifold with bounded geometry and $U \subset M$ is an open subset, then $U$ has totally bounded curvature, but may not have positive injectivity radius (so no bounded geometry).

## Boundary and bounded geometry

Let now $M_{0}$ be a Riemannian manifold with boundary, then $r_{\text {inj }}\left(M_{0}\right)=0$, so a manifold with non-empty boundary will never have bounded geometry in the sense of the above definition. The way around this conundrum was found by Schick (2001), who has defined the concept of "manifold with boundary and bounded geometry," (we shall call these manifolds "manifolds with boundary and relative bounded geometry," to avoid confusions). We recall the equivalent definition of manifolds with boundary and relative bounded geometry in (Ammann-Grosse-V.N.)

The main point of that definition is to assume that the boundary $\partial M_{0}$ of $M_{0}$ is a suitable submanifold of a (boundaryless) manifold $M$ with bounded geometry.

## Second fundamental form

Let hence $M$ be a (boundaryless) manifold with bounded geometry and let us consider a hypersurface $H \subset M$, that is, a submanifold $H$ of $M$ of codimension $\operatorname{dim}(M)-\operatorname{dim}(H)=1$. We assume that $H$ carries a globally defined unit normal vector field $\nu$. We let

$$
\exp ^{\perp}(x, t):=\exp _{x}^{M}\left(t \nu_{x}\right)
$$

be the exponential in the direction of the chosen unit normal vector. We shall need the second fundamental form $\mathrm{II}^{H}$ of $H$ in $M$, which, we recall, is defined by

$$
\Pi^{H}(X, Y) \nu:=\nabla_{X}^{M} Y-\nabla_{X}^{H} Y
$$

where $\nabla^{Z}$ is the Levi-Civita connection of $Z$. Equivalently, since $g\left(\nu, \nabla_{X}^{H} Y\right)=0$, we have $I^{H}(X, Y):=g\left(\nu, \nabla_{X}^{M} Y\right)$.

## Bounded geometry hypersurface

## Definition

Let $(M, g)$ be a Riemannian manifold of bounded geometry and $H \subset M$ be a hypersurface with unit normal vector field $\nu$ on $H$. We say that $H$ is a bounded geometry hypersurface in $M$ if:
(1) $H$ is a closed subset of $M$;
(1) all covariant derivatives $\left(\nabla^{H}\right)^{k} I^{H}, k \geq 0$, are bounded;
(T) $\exp ^{\perp}: H \times(-\delta, \delta) \rightarrow M$ is a diffeomorphism onto its image for some $\delta>0$.

Then $H$ has bounded geometry.
If $H \subset M$ is a compact hypersurface, then $H$ is a bounded geometry hypersurface in $M$.

## Manifolds with boundary and relative bounded geometry

We are ready now to recall the definition of a central concept in analysis on manifolds.

## Definition

We shall say that $M_{0}$ is a manifold with boundary and relative bounded geometry if $M_{0}$ is isometrically contained in a (boundaryless) Riemannian manifold $M$ with bounded geometry such that $\partial M_{0}$ is a bounded geometry hypersurface in $M$.

Note that we use the term "manifold with boundary and relative bounded geometry," which we think is more precise than the term "manifold with boundary and bounded geometry" used before.

## Trace and 'negative' spaces

Let $\partial_{\nu}$ be the normal derivative at the boundary and let

$$
H_{0}^{k}\left(M_{0}\right):=\cap_{j=0}^{k-1} \operatorname{ker}\left(\operatorname{res} \circ \partial_{\nu}^{j}\right)
$$

denote the joint kernel of the restrictions maps res $\circ \partial_{\nu}^{j}, 0 \leq j \leq k-1$.
Theorem (Trace theorem: Grosse-Schneider, Triebel)
Let $M_{0}$ be a manifold with boundary and relative bounded geometry. Then, for every $s>1 / 2$, the restriction res: $\mathcal{C}_{c}^{\infty}\left(M_{0}\right) \rightarrow \mathcal{C}_{c}^{\infty}\left(\partial M_{0}\right)$ extends by continuity to a surjective map

$$
\text { res: } H^{s}\left(M_{0}\right) \rightarrow H^{s-\frac{1}{2}}\left(\partial M_{0}\right)
$$

Moreover, $\mathcal{C}_{c}^{\infty}\left(M_{0} \backslash \partial M_{0}\right)$ is dense in $H_{0}^{k}\left(M_{0}\right)$ and $H^{-k}\left(M_{0}\right)$ identifies with $H_{0}^{k}\left(M_{0}\right)^{*}, k \in \mathbb{N}$.

## Finite width

## Definition

Let $\left(M_{0}, g\right)$ be a Riemannian manifold with boundary $\partial M_{0}$. We say that $M_{0}$ has finite width if:
(1) $\left(M_{0}, g\right)$ is a manifold with boundary and relative bounded geometry and
(1) The function $M_{0} \ni x \rightarrow \operatorname{dist}_{M_{0}}\left(x, \partial M_{0}\right)$ is bounded on $M_{0}$.

The last condition is equivalent to
" $\exists R>0$ such that $M_{0} \subset\left\{x \in M \mid \exists y \in \partial M_{0}, \operatorname{dist}_{M_{0}}(x, y)<R\right\}$."

## Poincaré inequality

## Theorem (Ammann-Grosse-V.N., Sakurai)

Let $\left(M_{0}, g\right)$ be a Riemannian manifold with finite width. Then there exists $0<C_{M_{0}}<\infty$ such that, for all $f \in \mathcal{C}_{c}^{\infty}\left(M_{0}\right)$ (thus $f=0$ on the boundary of $M_{0}$ ),

$$
\|f\|_{L^{2}\left(M_{0}\right)} \leq C_{M_{0}}\|d f\|_{L^{2}\left(M_{0}\right)} .
$$

A regularity argument then yields the following result.

## Theorem (Ammann-Grosse-V.N.)

Let $M_{0}$ be a smooth Riemannian manifold with smooth boundary $\partial M_{0}$ and finite width. Then $\Delta$ induces isomorphisms

$$
\Delta_{D}=\Delta: H^{m+1}\left(M_{0}\right) \cap H_{0}^{1}\left(M_{0}\right) \rightarrow H^{m-1}\left(M_{0}\right), \quad m \in \mathbb{Z}_{+}:=\{0,1, \ldots\}
$$

## Applications: polygonal regularity

Let us consider a polygonal domain $\Omega_{\infty}$ with maximum angle $\alpha_{M A X}$ and the Poisson problem

$$
\Delta u_{\infty}=f_{\infty} \text { in } \Omega_{\infty}, \quad u_{\infty}=0 \text { on } \partial \Omega_{\infty}
$$

It is known that $u_{\infty}$ has limited regularity, less than $H^{1+\pi / \alpha_{M A X}}$ : Costabel, Dauge, Griesvard, Kondratiev, Mazya, ...

## Theorem (Limited regularity: Kondratiev, ... )

Let $\delta<\pi / \alpha_{\text {MAX }}$. Then there exists $C=C_{\delta}>0$ such that the solution $u_{n} \in H_{0}^{1}\left(\Omega_{n}\right)$ of the equation $\Delta u_{n}=f_{n}$ satisfies

$$
\left\|u_{n}\right\|_{H^{1+\delta}\left(\Omega_{n}\right)} \leq C\left\|f_{n}\right\|_{L^{2}\left(\Omega_{n}\right)} .
$$

## Applications: well-posedness in weighted spaces

The "limited regularity" theorem of the previous slide is not very satisfying. It is, in any case, a consequence of the following more satisfying result in the Babuška-Kondratiev (weighted Sobolev) spaces

$$
\mathcal{K}_{a}^{m}(\Omega):=\left\{\rho^{|\alpha|-a} \partial^{\alpha} u \in L^{2}(\Omega),|\alpha| \leq m\right\}
$$

$\rho=$ distance to vertices (Costabel, Dauge, Kondratiev, Mazya, ... )

## Theorem (Unlimited weighted regularity: Kondratiev '67, ...)

Let $0 \leq \delta<\pi / \alpha_{\text {MAX }}$ and $m \in \mathbb{Z}_{+}$. Then there exists $C=C_{m, \delta}>0$ such that the solution $u_{\infty} \in H_{0}^{1}\left(\Omega_{\infty}\right)$ of the equation $\Delta u_{\infty}=f_{\infty}$ satisfies

$$
\left\|u_{\infty}\right\|_{\mathcal{K}_{\delta+1}^{m+1}\left(\Omega_{\infty}\right)} \leq C\left\|f_{\infty}\right\|_{\mathcal{K}_{\delta-1}^{m-1}\left(\Omega_{\infty}\right)} .
$$

This, as well as the next theorem, are consequences of the well-posedness on manifolds with boundary and relative bounded geometry (Thm 68 stated earlier).

## Applications: "rounding up the corners"

Consider a sequence of smooth domains $\Omega_{n}, n \in \mathbb{Z}$, converging to a polygonal domain $\Omega_{\infty}$ (picture) and the associated Poisson problems $\Delta u_{n}=f_{n}, u_{n} \in H_{0}^{1}\left(\Omega_{n}\right)$, for which we want uniform estimates.


## Applications: "rounding up the corners"

## Theorem (Daniel-Labrunie-V.N.)

Let $\Omega_{n} \subset \mathbb{R}^{2}$ be a sequence of smooth domains "converging" to a polygonal domain $\Omega_{\infty}$, as in the picture (repeated below). Let $\delta<\pi / \alpha_{M A X}$. Then there exists $C=C_{\delta}>0$ such that for all $n \in \mathbb{N} \cup\{\infty\}$, the solution $u_{n} \in H_{0}^{1}\left(\Omega_{n}\right)$ of $\Delta u_{n}=f_{n}$ satisfies

$$
\left\|u_{n}\right\|_{H^{1+\delta}\left(\Omega_{n}\right)} \leq C\left\|f_{n}\right\|_{L^{2}\left(\Omega_{n}\right)} .
$$



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## REMOVED STUFF 2

## Ellipticity

$M$ is still a closed manifold (smooth, compact, without boundary).

## Theorem

Let $a \in S^{m}\left(T^{*} M\right)$. The following are equivalent
(1) There exist $b \in S^{-m}\left(T^{*} M\right)$ such that $a b-1 \in S^{-1}\left(T^{*} M\right)$
(2) There are $C, R>0$ such that $|a(\xi)| \geq C|\xi|^{m}$ for $\xi \in T^{*} M,|\xi| \geq R$.

Then a is called elliptic and this property is a property of its class in $S^{m}\left(T^{*} M\right) / S^{m-1}\left(T^{*} M\right)$.

## Definition

$P \in \Psi^{m}(M)$ is elliptic if $\sigma_{m}(P) \in S^{m} / S^{m-1}\left(T^{*} M\right)$ is elliptic.
We can easily include vector bundles.

## Fredholm and Atkinson

Recall the following that if $X$ and $Y$ are Banach spaces, an operator $T \in \mathcal{B}(X ; Y)$ is Fredholm if

$$
\operatorname{dim} \operatorname{ker}(P), \operatorname{dim}(Y / P X)<\infty
$$

(It implies $P X=$ closed.)

## Theorem (Atkinson's theorem)

Let $X$ and $Y$ be Banach spaces; $T \in \mathcal{B}(X ; Y)$ is Fredholm if, and only if, there exists $Q \in \mathcal{B}(Y ; X)$ such that both $T Q-1_{Y}$ and $Q T-1_{X}$ are compact.

## Fredholm property

## Theorem

Assume $M$ is a closed manifold and $P \in \Psi^{m}(M ; E)$ is elliptic. Then, for any $s \in \mathbb{R}$,

$$
P: H^{s}(M ; E) \rightarrow H^{s-m}(M ; E) \quad \text { is Fredholm. }
$$

## Proof.

- $P$ is elliptic means that there is $b \in S^{-m}\left(T^{*} M\right)$ such that $\sigma_{m}(P) b-1 \in S^{-1}\left(T^{*} M\right)$.
- The surjectivity of $\sigma_{-m}$ and its multiplicativity give that there exists
$Q \in \Psi^{-m}(M)$ (parametrix) such that $P Q-1, Q P-1 \in \Psi^{-1}(M)$.
- Consequently, $P Q-1$ and $Q P-1$ are compact operators. Atkinson's theorem then gives that $P$ is Fredholm.

The converse is true (Fredholm for one s implies elliptic).

## Extensions of the Fredholmness results

- A quite satisfying result! However, much is hidden. The pseudodifferential operators that we defined are not enough, they contain too few regularizing operators!
- We need to enlarge the initial, small calculus (open pr. in gen).
- Open problem: find the "right class of regularizing operators."
- Second part of the lecture notes: general methods to extend some of the results in this talk (Pseudodifferential operators, Fredholm property, ...) to other classes of manifolds.


[^0]:    ${ }^{\text {a }}$ (Schwartz' kernel theorem, next!)

