



$$d(PMC, EZ) = 2$$

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# Table of Contents

1 Motivation

2 Waves

3 Heat (Joint work with Jiamin Zhu (LJLL-Paris6))

## So many reasons...

$$d(PMC, EZ) = 2$$

**Proof 1:**

$$d(PMC, EZ) = d(PMC, JV) + d(JV, EZ) = 1 + 1 = 2$$

**Proof 2:**

$$d(PMC, EZ) = d(PMC, VK) + d(VK, EZ) = 1 + 1 = 2$$

# Table of Contents

- 1 Motivation
- 2 Waves
- 3 Heat (Joint work with Jiamin Zhu (LJLL-Paris6))

ESAIM: Control, Optimisation and Calculus of Variations

April 1999, Vol. 4, p. 37–56

URL: <http://www.emath.fr/cocv/>

## WELL POSEDNESS AND CONTROL OF SEMILINEAR WAVE EQUATIONS WITH ITERATED LOGARITHMS \*

PIERMARCO CANNARSA<sup>1</sup>, VILMOS KOMORNIK<sup>2</sup> AND PAOLA LORETI<sup>3</sup>

**Abstract.** Motivated by a classical work of Erdős we give rather precise necessary and sufficient growth conditions on the nonlinearity in a semilinear wave equation in order to have global existence for all initial data. Then we improve some former exact controllability theorems of Imanuvilov and Zuazua.

**Résumé.** Motivé par un travail classique d'Erdős on donne des conditions nécessaires et suffisantes de croissance de la non linéarité dans une équation des ondes semilinéaire pour l'existence des solutions globales pour toutes les données initiales. Ensuite on améliore certains théorèmes antérieurs de contrôlabilité exacte de Imanuvilov et de Zuazua.

Roughly they establish the controllability of the  $1 - d$  semilinear wave equation for nonlinearities of the form

$$f(s) = s[\log[\log[\log \dots |s|]]]^2$$

**Proof**<sup>2</sup>

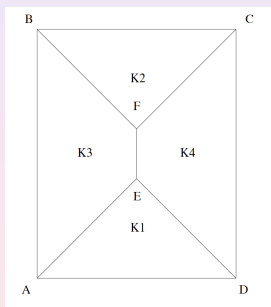


Figure 2 in their paper

<sup>2</sup>Similar sidewise energy arguments have been used in other contexts such as waves on networks: J. Lagnese, G. Leugering and G. Schmidt...

# Why?

The nonlinearity

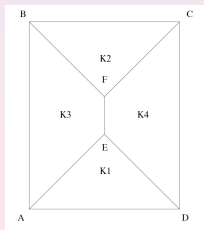
$$f(s) = s[\log[\log[\log \dots |s|]]]^2$$

is *subcritical* both for

$$y_{tt} - y_{xx} + f(y) = 0$$

and

$$y_{xx} - y_{tt} - f(y) = 0$$



# Osgood condition

For first order ODEs

$$x'(t) = f(x(t))$$

the criticality condition is

$$\int_{z_0}^{\infty} \frac{1}{f(z)} dz = \infty.$$

This condition ensures global well-posedness, while finite-time blow up occurs if the integral converges.

Thus

$$f(s) = s[\log[\log[\log \dots |s|]]]$$

is **critical**.



## 2nd order Osgood condition

For second order ODEs

$$x''(t) = f(x(t))$$

multiplying by  $x'(t)$  we get

$$\frac{1}{2}|x'(t)|^2 - F(x(t)) = C$$

And therefore

$$x'(t) = \pm \sqrt{2(C + F(x(t)))},$$

the critical threshold becomes

$$f(s) = s[\log[\log[\log \dots |s|]]]^2.$$

## Multi-d

**The multi-d analog of this result is unknown**, because of the ill-posedness of the sidewise wave equation.

A partial result is known for non-linearities of the order of <sup>3</sup>

$$|f(s)| \ll |s| \log^{1/2}(|s|), \text{ as } |s| \rightarrow \infty.$$

Employing techniques based on Carleman inequalities one could expect to extend this result to the range

$$|f(s)| \ll |s| \log^{3/2}(|s|), \text{ as } |s| \rightarrow \infty,$$

as it occurs for the heat equation.

But getting to the 1-d range seems hard although nothing excludes the result to be true, as far as we know.

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<sup>3</sup>X. Zhang and E. Z. Exact controllability of the semi-linear wave equation. In “Unsolved problems in mathematical systems and control theory”, Princeton University Press, 2004, pp. 173-178.

This  $2/3$  exponent may not be improved through Carleman estimates as pointed out by Th. Duyckaerts, X. Zhang and EZ (Annales IHP, 2005), based on the following result by V. Z. Meshkov, 1991.

### Theorem

(Meshkov, 1991). Assume that the space dimension is  $n = 2$ . Then, there exists a nonzero complex-valued bounded potential  $q = q(x)$  and a non-trivial complex valued solution  $u = u(x)$  of

$$\Delta u = q(x)u, \quad \text{in } \mathbb{R}^2, \quad (1)$$

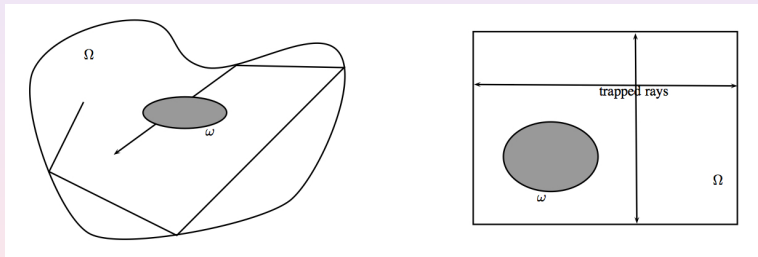
with the property that

$$|u(x)| \leq C \exp(-|x|^{4/3}), \quad \forall x \in \mathbb{R}^2 \quad (2)$$

for some positive constant  $C > 0$ .

## What if the control time is too short?

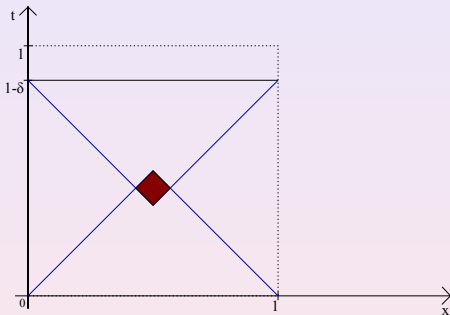
These results hold for large enough time horizons of control. Roughly, under the so-called Geometric Control Condition (GCC) by Bardos - Lebeau - Rauch (1988). It asserts, roughly, that all rays of geometric optics enter the control set  $\omega$  in time  $T$ .



When this condition holds all data can be controlled.  
But what when  $T$  is too short for GCC to be fulfilled?

1-d :  $T < 2$ 

In 1-d if  $T < 2$  the space of controllable initial data can be characterised either using characteristic arguments or Fourier series expansions:<sup>4 5</sup>



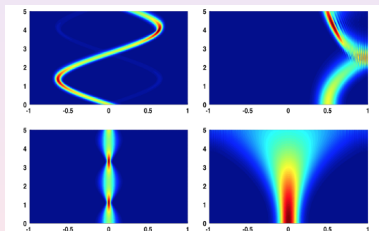
<sup>4</sup>J. Lohéac and E. Zuazua, Norm saturating property of time optimal controls for wave-type equations, CPDE 2016, Bertinoro, 2016, IFAC-PapersOnLine, 49 (8) 37-42, 2016.

<sup>5</sup>D. A. Ivanov, M. M. Potapov. Approximations to Time-Optimal Boundary Controls for Weak Generalized Solutions of the Wave Equation. Computational Mathematics and Mathematical Physics, 2017, Vol. 57, No. 4, pp. 607-625.

# Numerical algorithms

How to build numerical algorithms enabling to compute the minimal control time for a specific initial datum?

Note that, so far, all effort have been devoted to build numerical algorithms enabling to get the GCC. But nothing has been done for shorter  $T$ 's.<sup>6, 7</sup>



<sup>6</sup>E. Z. Propagation, observation, and control of waves approximated by finite difference methods. SIAM Review, 47 (2) (2005), 197-243.

<sup>7</sup>A. Marica y E. Zuazua, Propagation of 1 –  $D$  waves in regular discrete heterogeneous media: A Wigner measure approach, Found. Comput. Math. (2015) 15:1571-1636.

# Table of Contents

- 1 Motivation
- 2 Waves
- 3 Heat (Joint work with Jiamin Zhu (LJLL-Paris6))

## Motivation

Beyond standard control problems in which the control enters as an external applied source, many problems arising in population dynamics, ecology, etc raise interesting questions about the possibility of controlling the dynamics regulating some of the parameters entering in the system.<sup>8,9</sup>

Approximate controllability of

$$y_t - y_{xx} + v(x, t)y + f(y) = 0,$$

employing intermitent alternating controls: From initial to final data with the same number of zeroes...

- ① Phase 1:  $v \equiv 0$ , free evolution of the semilinear equation.
- ② Phase 2:  $v \equiv v(x)$ , static control.

<sup>8</sup>P. M. Cannarsa, G. Floridia, A.Y. Khapalov, Multiplicative controllability for semilinear reaction-diffusion equations with finitely many changes of sign, arxiv 2015, JMPA, to appear.

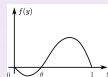
<sup>9</sup>Motivated by tumor growth, fissionable materials,...



# Numerical control of travelling wave solutions: Jiamin Zhu (LJLL-Paris6) and E. Z.

The initial value problem

$$y_t - y_{xx} = ay(1 - y)(y - \theta), \quad y(0, \cdot) = y_0, \quad x \in \mathbf{R}, \quad t \in \mathbf{R}^+,$$

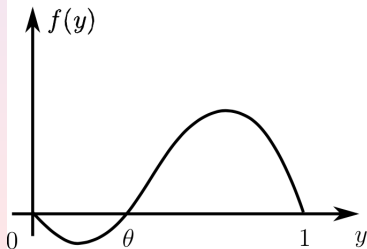


- Represents the net population change from birth and death.
- Typical application : spread of invading organisms in ecology systems, (cf. e.g. M.A. Lewis and P. Kareiva 1993).
- The role of parameters  $a, \theta$ :
  - $a > 0$  : reproductive rate
  - $\theta \in (0, 1)$  : local critical density or Allee threshold that determines the sign (positive or negative) the population growth.
- Other applications : population genetics (biology), propagation of nerve tension (neurobiology), propagation waves in chemical reactors (chemistry), etc.

## Our original motivation: multilingualism

# A Game-Theoretic Analysis of Minority Language Use in Multilingual Societies<sup>1</sup>

José-Ramón Uriarte  
University of the Basque Country  
October 2015

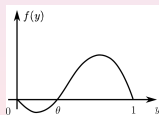


**Allee effet** : population growth is negative (leading to the extinction) when the density of the population is lower than the Allee threshold  $\theta$ , otherwise the population will reach carrying capacity.

- increase  $\theta$  by the sterile male technique, the mating disruption (pest management technique), etc.
- decrease  $\theta$  by providing protection (e.g. efficient feeding, suppressing natural enemies) to the population

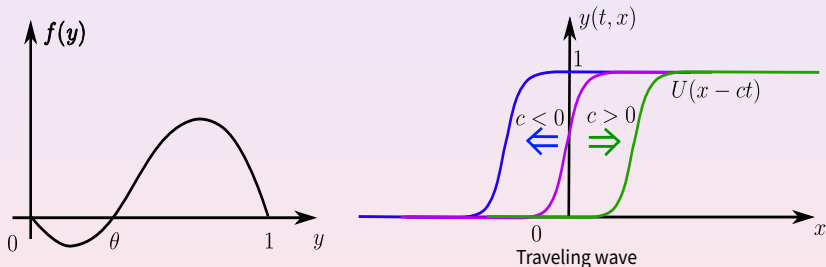
**Typical solutions of the system** :

- Steady state constant solutions :  $y \equiv 0$  or  $\theta$  or 1.
- Traveling wave solutions : link two of the three constant solutions



# Traveling wave solutions

- solution of the form  $y(t, x) = U(x - ct)$ ,  $U(\pm\infty) = U_{\pm}$ ,  
 $U(\pm\infty)' = 0$ ,  
 where  $U(x)$  is the wave profile and  $c$  is the wave speed.
- sign of the wave speed :  $\text{sign } c = - \int_0^1 f(t) dt$



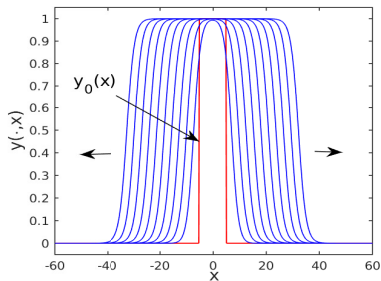
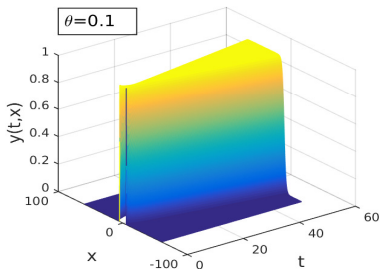
- Effect of control  $\theta$ 
  - change the wave speed
  - when  $U_- = 0$ ,  $U_+ = 1$ ,  $c = \sqrt{a/2}(2\theta - 1)$  and  $U(x) = \frac{e^{\sqrt{a/2}x}}{1 + e^{\sqrt{a/2}x}}$

# La ola



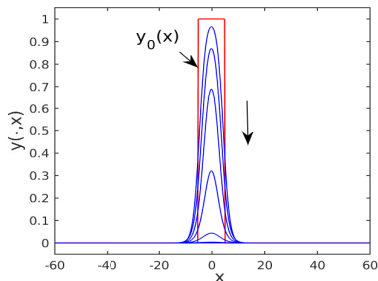
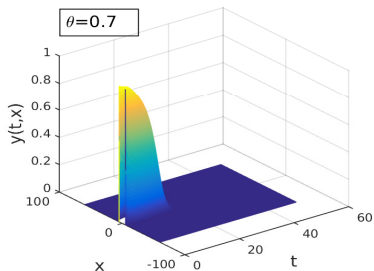
# Is it possible to control the traveling waves ?

- $\theta = 0.1$  : invasion of the population



# Is it possible to control the traveling waves ?

- $\theta = 0.1$  : invasion of the population
- $\theta = 0.7$  : extinction of the population



# Formulation of the control problem

## Control problem $\mathcal{P}_c$

Find  $\theta(t) \in [0, 1]$ ,  $t \in [0, T]$  such that the solution of

$$y_t - y_{xx} = ay(1 - y)(y - \theta),$$

$$y(0, \cdot) = y_0(\cdot)$$

develops into an expected wave  $U(\cdot)$  at the given time  $T$ , minimizing

$$J(\theta) = \left| y(T, \cdot) - U \right|^2$$



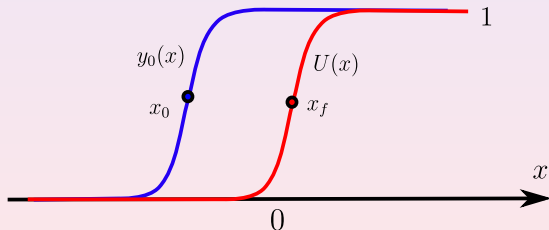
# Explicit solutions : from traveling wave to traveling wave

- $U(x - c_d t)$  is a traveling wave : given  $\theta_d \in (0, 1)$ ,  $x_d \in \mathbf{R}$

$$U(x) = e^{\sqrt{a/2}(x-x_f)} / (1 + e^{\sqrt{a/2}(x-x_f)}),$$

$$c_d = \sqrt{a/2}(2\theta_d - 1)$$

- $y_0(\cdot)$  is a translation of  $U(\cdot)$ , i.e.,  $y_0 = U(x + x_f - x_0)$



Wave speed  $c$  is bounded  $\Rightarrow T$  must be sufficiently large, i.e.,  $T \geq T_{min}$

# Piecewise constant control strategy (infinite many choices)

$$\theta^*(t) = \begin{cases} \theta^1, & t \in (0, t_1), \\ \theta^2, & t \in [t_1, t_1 + h), \\ \dots & \\ \theta^N, & t \in [T - h, T], \end{cases}$$

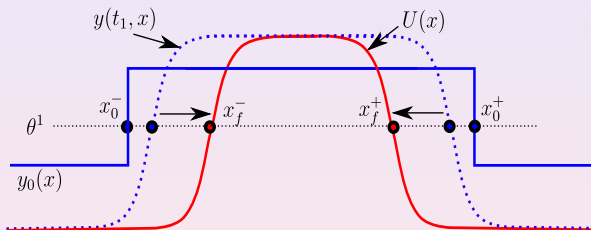
Choose  $\theta^i$ ,  $i = 2, \dots, N$  such that  $h \sum_{i=2}^N c(\theta^i) = (x_f - x'_0 - c_1 t_1)$ , where  $c(\theta^i) = \sqrt{a/2}(2\theta^i - 1)$

Since  $c(\theta^i) < \sqrt{|a_1/2|} \Rightarrow T$  should be sufficiently large, i.e.,

$$T > T_{min} = t_1 + \frac{|x_f - x'_0 - c_1 t_1|}{\sqrt{|a_1/2|}}$$

## Quasi explicit examples

One can use the existing results on the stability and the attractiveness of the dynamics for  $\theta = \text{constant}$  fixed to "approximately" control the system in two steps.<sup>10, 11</sup>



<sup>10</sup>D.G. Aronson, H.F. Weinberger, Nonlinear diffusion in population genetics, combustion, and nerve pulse propagation, in PDE and Related Topics, Lecture Notes in Math., vol. 446, Springer, Berlin, 1975, pp. 5-49.

<sup>11</sup>P.C. Fife, J.B. McLeod, The approach of solutions of nonlinear diffusion equations to travelling front solutions, ARMA, 1977, vol. 65 no. 4, p. 335-361.

## Numerics: Gradient descent

- Start with an initialization  $u^0 = (\theta^1, \dots, \theta^N, t_1)$ .
- At step  $k$  :
  - step 1 : integrate the state system with control  $u^k$  from  $t = 0$ ,  $y(0, \cdot) = y_0$  to  $t = T \Rightarrow$  output  $y(t, x)$
  - step 2 : compute the cost  $J(u^k)$ . If  $J(u^k) < \epsilon$ , stop.
  - step 3 : integrate the adjoint system from  $t = T$ ,  $p(T, \cdot) = 2(\bar{y}(T, \cdot) - Y_d)$  to  $t = 0 \Rightarrow$  output  $p(t, x)$
  - step 4 : calculate the gradient of the cost functional by  $J'(u^k) = \int_0^T \int_{\Omega} p(t, x) \partial_u f(y, u^k) dx dt$ , where  $p$  is the adjoint state.
  - step 5 : set the descend direction  $d_k = J(u^k)$ , compute a descend step size  $s_k$ , and set the new control by  $u^{k+1} = u^k + s_k d_k$ . Return to step 1.
- non convex optimal control problem  $\Rightarrow$  numerically obtained optimal controls depend on the initialization  $u^0$

# Two-grids

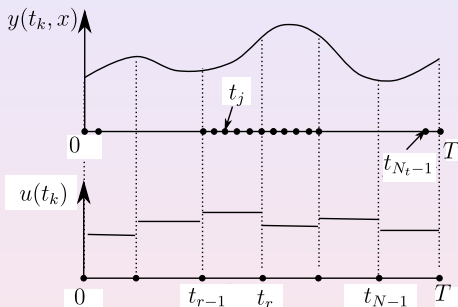
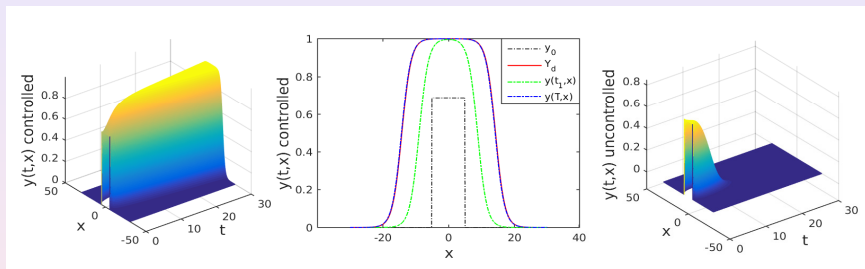


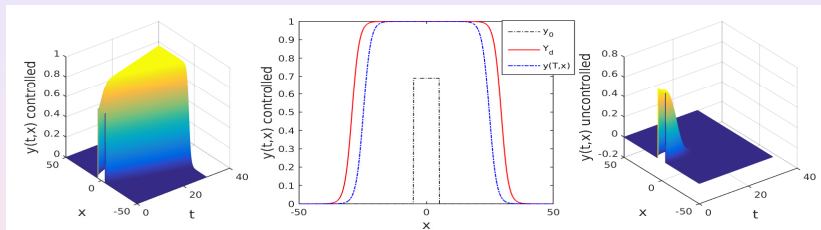
Figure: Space subdivisions for the state  $y$  and for the control  $u$  at time grid  $t_k$ .

# Numerical experiments: Fighting against extinction, $T > T_{min}$



**Figure:** Left and middle : solution of problem  $(\mathcal{P})_{opt}$  after control;  
Right : state  $y$  of uncontrolled system.

# Numerical experiments: Fighting against extinction, $T < T_{min}$



**Figure:** Left and middle : solution of problem  $(\mathcal{P})_{opt}$  after control;  
Right : state  $y$  of uncontrolled system.

# Perspectives

- So far numerical simulations: Theoretical analysis to be done beyond quasi-explicit examples.
- More complex models (multi-d, systems,...)



# Happy birthday, Piermarco. We will continue the celebration in Benasque August 20th...



**VII Partial differential equations, optimal design and numerics**

Centro de Ciencias de Benasque Pedro Pascual, Huesca-Spain  
2017, Aug 20 – Sep 01

**Organizers**  
G. Buttazzo (U. Pisa)  
O. Glass (Université Paris-Dauphine)  
G. Leugering (U. Erlangen-Nürnberg)  
E. Zuazua (DeustoTech - UAM - LJLL (Paris))

The workshop is intended to provide a fruitful atmosphere for discussions and joint research work on themes involving partial differential equations and their applications to shape optimization, optimal control problems, singularities in fracture mechanics and fluid dynamics, and numerical analysis.

The work will be focused on new research trends in the fields above, in order to stimulate collaborations among participants by means of various activities on the basis of a daily programme (talks, seminars, minicourses, discussions...)

The activity is mainly intended for young scientists as PhD students and post-docs, and will be held thanks to the participation of a number of world leader mathematicians.

Limited partial financial support will be available for young researchers.

\* Further information about registration, accommodation, travel can be found at <http://www.benasque.org/2017.pdf>

\* For further local information about Benasque and the region: <http://www.benasque.com>

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